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On the Gas Surrounding High-Redshift Galaxy Clusters*

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Abstract: Francis & Hewett (1993) identified two 10 Mpc-scale regions of the high-redshift universe that were seemingly very overdense in neutral hydrogen. Subsequent observations showed that at least one of these gas-rich regions enveloped a cluster of galaxies at redshift 2.38. We present improved observations of the three background QSOs with sightlines passing within a few Mpc of this cluster of galaxies. All three QSOs show strong neutral hydrogen absorption at the cluster redshift, suggesting that this cluster (and perhaps all high-redshift clusters) may be surrounded by a ~ 5 Mpc-scale region containing $\sim 10^{12} M_{\odot}$ of neutral gas. We show that if most high-redshift clusters are surrounded by such regions, the gas must be in the form of many small (< 1 kpc), dense ($> 0.03 \text{ cm}^{-3}$) clouds, each of mass $< 10^6 M_{\odot}$. These clouds are themselves probably gathered into > 20 kpc-sized clumps, which may be galaxy halos or protogalaxies. If this gas exists, it will be partially photoionised by the UV background. We predict the diffuse Ly α flux from this photoionisation, and place observational limits on its intensity.

Keywords: galaxies: clusters: individual (2142–4420) — galaxies: distances and redshifts — quasars: absorption lines

1 Introduction

What did galaxy clusters look like ten billion years ago? Simulations (e.g. Brainerd & Villumsen 1994; Jenkins et al. 1998; Cen 1998) suggest that the ancestors of present-day rich galaxy clusters contained very little mass at redshifts above 2: there simply hadn't been time enough to assemble much dark or baryonic matter. Even the richest proto-clusters would only have had about twice the average density of the universe. These proto-clusters would typically have been filamentary in shape, and, due to their small masses, would not have been in virial equilibrium.

Despite their modest overdensities, these proto-clusters could have been the sites of the first galaxy formation. If galaxies had formed in these proto-clusters but not in the field, the overdensity of galaxies within the proto-cluster would be much greater than the overdensity of mass (biasing, e.g. Fry 1996; Bagla 1998; Baugh et al. 1999; Tegmark & Peebles 1998).

The observational situation is less clear. Galaxy clusters are now routinely being studied out to redshifts $z \sim 1$. These clusters are remarkably similar to low-redshift clusters: they are massive virially bound objects with strong X-ray emission (e.g. Rosati et al. 1998; Deltorn et al. 1997). Thus galaxy clusters appear to be well established by redshift 1 (see also Renzini 1997; Donahue & Voit 1999).

Observations at higher redshifts still are sparse. It is now clear that galaxies at $z > 2$ are at least as strongly clustered as galaxies today (e.g. Heisler, Hogan & White 1989; Quashnock, Vanden Berk & York 1996; Malkan, Teplitz & McLean 1996; Steidel et al. 1998; Giavalisco et al. 1998; Campos et al. 1999; Pascarelle, Windhorst & Keel 1998; Adelberger et al. 1998; Djorgovski et al. 1999), as predicted by the biasing model. The structure, mass and galaxy populations of these galaxy concentrations are, however, obscure.

A growing body of evidence suggests that high-redshift clusters contain substantial quantities of gas. The strongest evidence comes from high-redshift radio galaxies, which are believed to sit in dense cluster environments (e.g. Pentericci et al. 1997; Ivison et al. 2000). Many show high rotation measures (Carilli et al. 1997; Pentericci et al. 2000), possibly extended X-ray emission (Carilli et al. 1998), extensive emission-line halos (Bicknell et al. 2000) and associated absorption lines (Binette et al. 2000), suggesting the presence of dense inhomogeneous gas around the galaxies.

Are the clusters around radio galaxies typical? The evidence for gas in other high-redshift clusters is sparse. Francis et al. (1996) and Steidel et al. (1998) both found QSO absorption lines seemingly coincident with high- z clusters. Tentative Sunyaev–Zel'dovich effect decrement measurements have been claimed around other possible high- z clusters (e.g. Campos et al. 1999). Diffuse 'blobs' of Ly α emission on the 100 kpc scale have been found in the centres of two high- z clusters (Francis et al. 1996; Francis, Woodgate & Danks 1997; Steidel et al. 2000).

* Based on observations carried out at the Anglo Australian Telescope, Cerro Tololo Interamerican Observatory and Siding Spring Observatory.

Most of these observations suggest the presence of hot ionised gas in the central few hundred kpc of the clusters. Francis & Hewett (1993), however, presented tentative evidence for the existence of large concentrations of neutral gas extending over Mpc scales in the high-redshift universe. One of their concentrations was associated with a group of galaxies (Francis, Woodgate & Danks 1997). Their evidence was not conclusive, and there are currently no such data available for other clusters, but it is at least possible that similar extended halos of neutral gas surround all high-redshift clusters.

In this paper, we concentrate on one of Francis & Hewett's clusters: the 2142–4420 galaxy concentration at redshift 2.38. In Section 2 we present new data on this cluster, strengthening the case for a large halo of neutral gas surrounding it. In Section 3 we investigate ways of reconciling high neutral gas masses with cosmological models. In Section 4 we describe our attempt to observe diffuse Ly α emission from the hypothesised gas halo, and place upper limits on its intensity. Except where stated, we assume $H_0 = 70 \text{ km s}^{-1} \text{ Mpc}^{-1}$, and an open universe with $\Omega_0 = 0.2$ and $\Lambda =$ either 0 or 0.8 (both values of Λ give the same angular size distance and luminosity distance at redshift 2.38).

2 Observations of the 2142–4420 Cluster

2.1 Review of the 2142–4420 Cluster Properties

The 2142–4420 cluster lies at redshift 2.38, at coordinates 21:42:30–44:20:30 (J2000). As will be shown, it is a region of the early universe that is highly overdense in Ly α emitting galaxies. It is unlikely to be gravitationally bound, and so would not meet most low-redshift definitions of a galaxy cluster.

Francis & Hewett originally identified the cluster as a pair of strong Lyman-limit systems at matching redshifts (2.38) in two 19th magnitude QSOs: 2138–4427 and 2139–4434. The QSO sightlines are $500''$ (~ 4 proper Mpc) apart at $z = 2.38$ (both QSOs lie at $z \sim 3.2$). Such a pair of Lyman-limit systems at matching wavelengths is unlikely to occur by chance in Francis & Hewett's sample. As we report in Section 2.2, a third QSO has been found behind the cluster, and it too shows strong Ly α absorption at the cluster redshift.

Are there any galaxies associated with this concentration of QSO absorption lines? Many high-redshift galaxies show *weak* Ly α emission (Steidel et al. 1996; Hu, Cowie & McMahon 1998). The field containing the three QSOs, however, contains three *strong* Ly α -emitting galaxies at $z = 2.38$, with Ly α fluxes $F_{\text{Ly}\alpha} > 10^{-16} \text{ erg cm}^{-2} \text{ s}^{-1}$ ($> 10^{-19} \text{ W m}^{-2}$; Francis et al. 1996; Francis, Woodgate & Danks 1997), an order of magnitude greater than those of normal galaxies at these redshifts. All three have reliable spectroscopic redshifts.

Do three such sources constitute a cluster? Strong Ly α emitters are rare at these redshifts: Francis et al. surveyed a total volume of 460 co-moving cubic Mpc, but all three sources were found within a 5 cubic co-moving

Mpc volume, i.e. $\sim 1\%$ of the surveyed co-moving volume. Martínez-González et al. (1995) surveyed a co-moving volume of 1400 cubic Mpc at $z = 3.4$ for Ly α -emitting sources to a comparable flux limit but detected nothing. The odds of the cluster being an artifact of the coincidental proximity of three such sources is thus less than $(5/1860)^2$ (the probability of finding two more such objects within one proper Mpc of the first), i.e. $< 10^{-5}$. Clearly the space density of Ly α -emitting galaxies in this region is higher than average.

Could the presence of absorption in all three QSOs be coincidental, or is the cluster really surrounded by Lyman-limit absorption-line systems? The transverse separation of the QSO sightlines is ~ 5 Mpc, which corresponds to a redshift difference of $\sim 20 \text{ \AA}$ along the line of sight. The absorption-line systems all have equivalent widths of $> 20 \text{ \AA}$. The probability of seeing an absorption line with an equivalent width this strong within $\pm 20 \text{ \AA}$ of any given wavelength is 1.3% (Francis & Hewett 1993). Thus the probability of finding three such absorption-line systems within the cluster by chance is 2×10^{-6} .

This calculation should be regarded with caution: this region was first identified as interesting because of the absorption in the two original QSOs (Francis & Hewett 1993), so the statistics are *a posteriori*. The third QSO, however, was not involved in selecting this region for study, and its coincident absorption alone makes this region overdense in Lyman-limit systems with 98% confidence.

Figure 1 should make the geometry of the cluster clearer. Due to the low predicted overdensities of clusters at this redshift (Section 3.2), peculiar motions should be very small, so three-dimensional positions are plotted assuming that all redshift differences are due to distance. The three Ly α galaxies lie within 1 Mpc of each other. The absorption-line systems are far more dispersed, extending both to lower redshifts and transversely by ~ 5 Mpc.

The spatial extent and overdensity of this cluster are comparable to those of the clusters of Lyman-break galaxies being found by Steidel et al. (1998) at $3 < z < 3.5$: this cluster may be a representative of the same class of object.

2.2 Observations

Véron & Hawkins (1995) searched an area including this cluster for variable sources. In addition to both previously identified QSOs, they discovered a third QSO lying between the two: QSO 2139–4433 at $z = 3.22$ (i.e. at the same redshift as the other two background QSOs). We measured a position for QSO 2139–4433 (21:42:22.16–44:19:28.7, J2000) using our R-band image with an astrometric solution bootstrapped from online scans of UK Schmidt plates (Drinkwater, Barnes & Ellison 1995). A spectrum was obtained with the Low Dispersion Survey Spectrograph (LDSS, Colless et al. 1990) on the Anglo-Australian Telescope on the nights of 1996 August 13 and 14. The total exposure time was 47,700 s, and the spectral resolution 700 km s^{-1} . Part of the spectrum is shown in Figure 2.

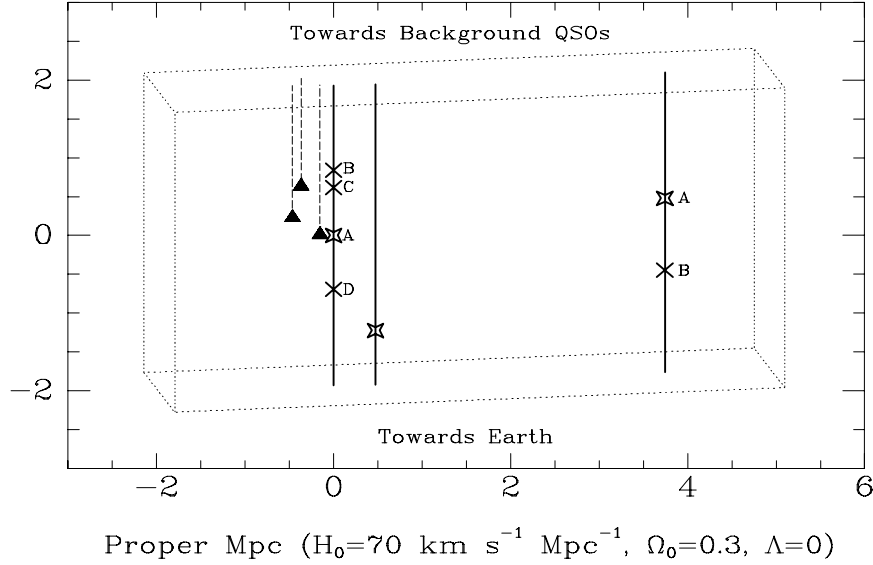


Figure 1 3D view of the 2142–4420 cluster. All objects have been assigned three-dimensional positions, based on their locations on the sky and their redshifts, assuming that all redshifts trace the Hubble flow. Solid lines are the sightlines to the three background QSOs. Stars are Ly-limit absorption systems, crosses are the lower-column-density absorption systems, and solid triangles are the Ly α -emitting galaxies. Redshifts increase upwards: galaxy redshifts have been derived, where possible, from the metal emission lines. From left to right, the QSO sightlines are those to 2139–4434, 2139–4433 and 2138–4427.

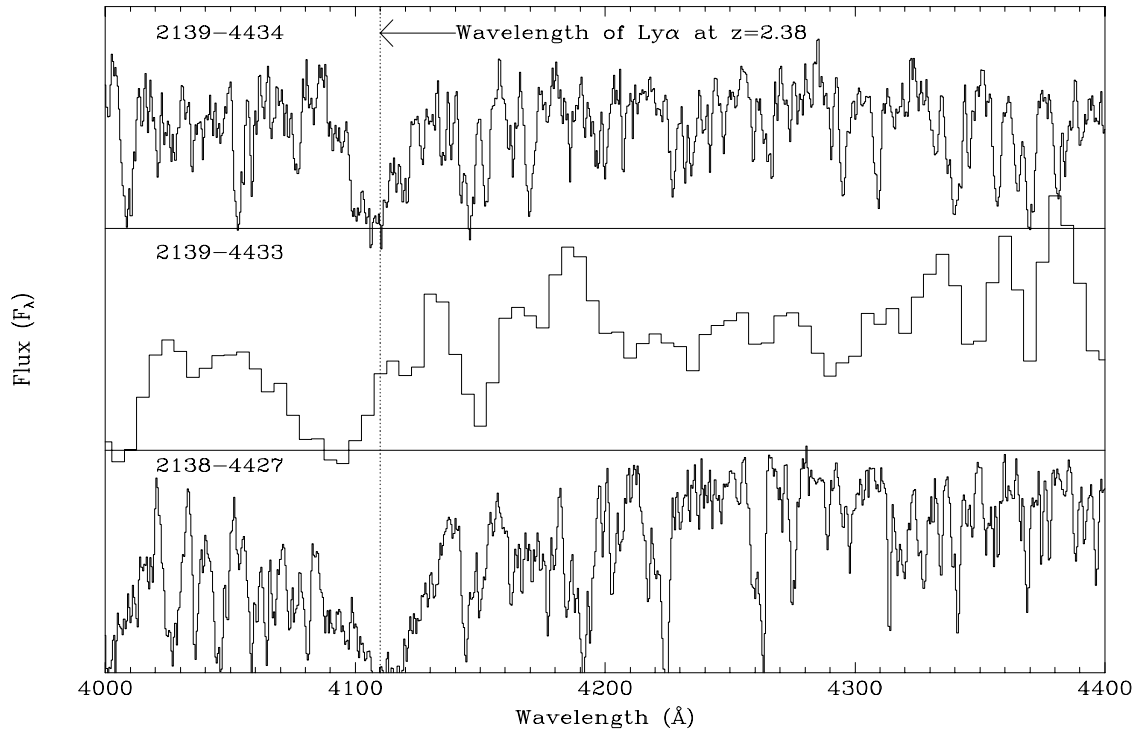


Figure 2 Spectra of the three background QSOs, showing the Ly α absorption at the cluster redshift (4108 Å). The top and bottom spectra are the original AAT spectra; the middle panel is the new LDSS spectrum.

As Figure 2 shows, the new QSO 2139–4433 has a strong absorption-line system close in wavelength to the absorption in the two previously known QSOs at $z = 2.38$. This further confirms the remarkable gas properties of this cluster.

Our original spectra of QSOs LBQS 2138–4427 and 2139–4434 are described by Francis & Hewett (1993).

Their resolution was excellent (full width at half maximum height 100 km s^{-1}) but the wavelength coverage ($4000\text{--}4600 \text{ Å}$) was small. An additional spectrum of QSO 2139–4434 was obtained with the KPGL1 grating in the Blue Air camera of the RC spectrograph on the CTIO 4 m telescope on 1995 August 20. Total exposure time was 12,000 s, with a spectral resolution of 200 km s^{-1} . This

spectrum, while inferior in resolution to the spectrum of Francis & Hewett, covers 3200–6200 Å: this greater wavelength range allows us to study CIV and Ly β absorption from the cluster.

2.3 Absorption Line Measurements

Combining the old and new data on the two brighter QSOs, we fit Voigt profiles interactively to the absorption at the cluster redshift, using the XVOIGHT program

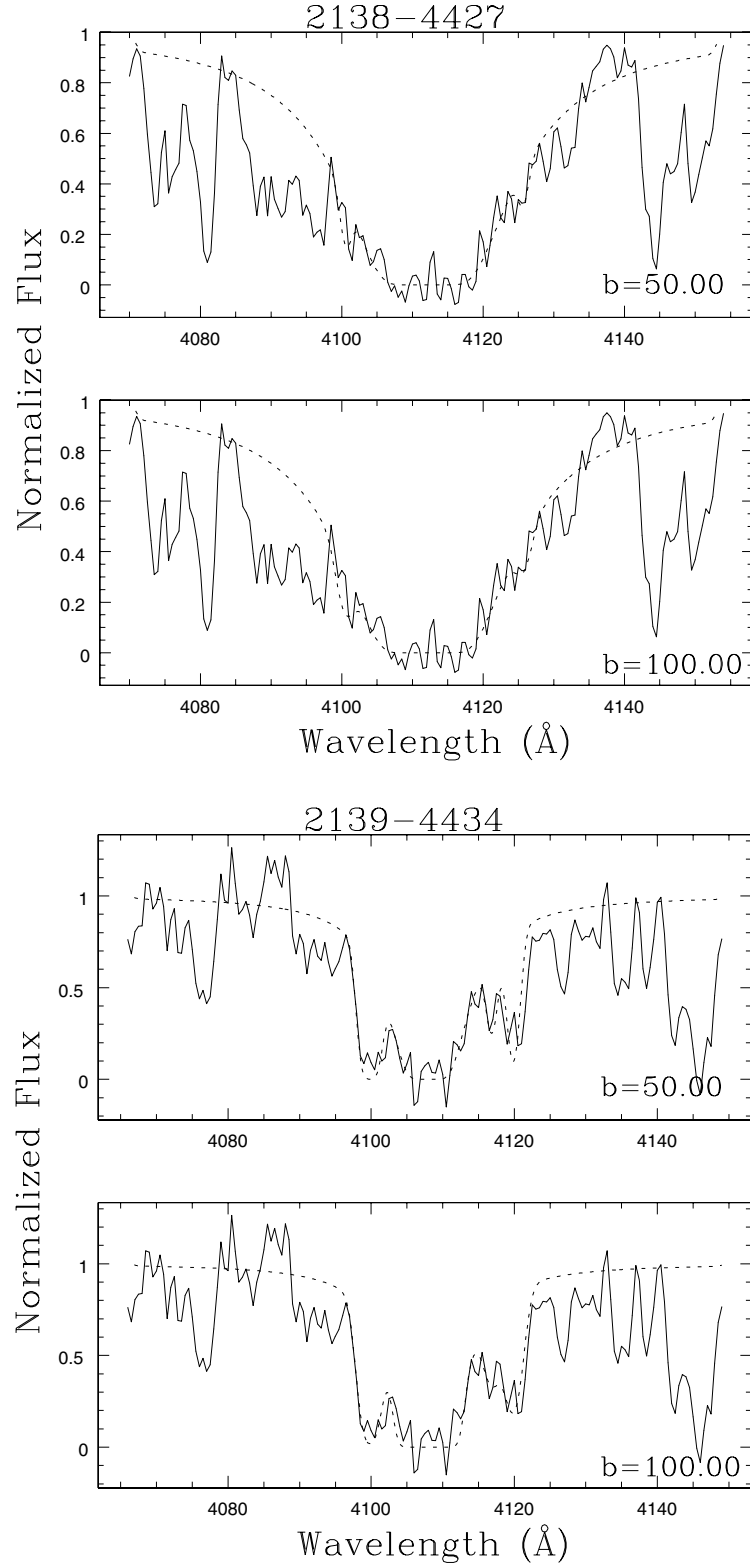


Figure 3 Voigt profile fits to the Ly α absorption at the cluster redshift in QSOs 2138–4427 and 2139–4434. Fits are shown for the two different velocity dispersions b assumed.

Table 1. Possible absorption systems in QSO 2138–4427

System	Ion	Redshift	Column density $\log(N[\text{cm}^{-2}])$	
			$b = 50 \text{ km s}^{-1}$	$b = 100 \text{ km s}^{-1}$
A	H I 121-6	2.3825	20.47	20.47
	Si II 130-4	2.3823	14.63	14.72
	Si II 126-0	2.3825	14.63	14.22
	Si III 120-6	2.3824	14.30	13.98
	O I 130-2	2.3820	16.22	15.47
	C I 127-7	2.3821	14.62	14.70
	C II 133-4	2.3822	16.00	15.13
B	H I 121-6	2.3731	13.97	14.10

Table 2. Possible absorption systems in QSO 2139–4434

System	Ion	Redshift	Column density $\log(N[\text{cm}^{-2}])$	
			$b = 50 \text{ km s}^{-1}$	$b = 100 \text{ km s}^{-1}$
A	H I 121-6	2.3792	19.80	19.67
	Si II 130-4	2.3794	13.63	13.80
	Si II 126-0	2.3792	13.53	13.67
	Si III 120-6	2.3789	13.45	13.55
	O I 130-2	2.3804	16.00	15.47
	C I 127-7	2.3790	14.00	14.17
	C II 133-4	2.3783	16.03	15.13
	C IV 154-8	2.3787	14.30	14.23
	C IV 155-1	2.3787	14.17	14.43
B	H I 121-6	2.3890	14.56	14.37
C	H I 121-6	2.3865	14.07	14.07
D	H I 121-6	2.3724	16.93	14.80

(Mar & Bailey 1995). The low spectral resolution, restricted wavelength coverage and blending in our spectra make this process a difficult and ambiguous one. Nonetheless, certain definite conclusions can be reached. Multiple components are required to obtain adequate fits to the Ly α absorption (Figure 3). A minimum of 2–4 components are required (Tables 1 and 2); many more, each with smaller column densities, give equally good fits. The column densities of the subsidiary systems are not well constrained. We could not determine the velocity dispersion b of the metal lines; an upper limit of $\sim 100 \text{ km s}^{-1}$ can be placed. The central component of the Ly α absorption in all three QSOs was broader: the flux touches zero over $\sim 100 \text{ km s}^{-1}$ or more.

We searched for metal-line absorption at the redshift of the dominant Ly α absorption components. With the exception of C IV, these all lie within the Ly α forest, and hence may be chance coincidences with the forest lines. The strongest line near the expected wavelength was fitted, assuming velocity widths of 50 and 100 km s^{-1} , and the results are shown in Tables 1 and 2. Due to the risk of blending or confusion with Ly α forest lines, the metal-line column densities should be taken as upper limits. Plots of the metal lines within the forest can be found in Francis & Hewett (1993). Note that in QSO 2138–4427, strong absorption lines were invariably detected at the expected

wavelengths, while in QSO 2139–4434 the lines were weaker and at slightly shifted wavelengths. We conclude that the central absorption component in QSO 2138–4427 does contain metals, roughly as measured, while for QSO 2139–4434 some or all of the putative lines (except C IV) may be misidentified Ly α forest lines. Our spectrum of QSO 2139–4433 had too low a resolution to determine anything other than the Ly α redshift (2.366) and column density ($\log(N_{\text{H}}) \sim 20.7$).

Are the absorption-line systems really Lyman-limit systems, or could they just be clusters of lower-column-density Ly α forest lines? In QSO 2138–4427 the Ly α line shows broad wings, and strong absorption is seen at the expected wavelength of most common metal absorption lines; it therefore seems probable that this is, as modelled, a high-column-density absorption system, probably lying on the column density borderline between Lyman-limit and damped Ly α systems. The spectrum of QSO 2139–4433 is of too low resolution to say much, but the great width and equivalent width of the Ly α absorption also suggest that its absorption column is large.

In QSO 2139–4434, however, the situation is more ambiguous. It is possible to fit the Ly α absorption either with a single absorption-line system with column density $N_{\text{H}} \sim 10^{19.7} \text{ cm}^{-2}$ (plus three much weaker components in the wings), or with a blend of weaker Ly α lines, spread

over $\sim 200 \text{ km s}^{-1}$ and with a combined neutral hydrogen column density that can be as low as $N_{\text{H}} \sim 10^{16.5} \text{ cm}^{-2}$. Two pieces of evidence support this latter fit. First, the redshifts of the supposed metal lines vary by $\pm 80 \text{ km s}^{-1}$ (although some or all may be chance coincidences with Ly α forest lines). This can be explained if they are coming from different subcomponents of the absorption system. Second, there is tentative evidence that Ly β absorption is weak: the spectrum is poor at this wavelength, and the continuum hard to define, but the Ly β absorption can be well fitted with column densities as low as $N_{\text{H}} \sim 10^{16}$ (although much greater columns also give acceptable fits). On the other hand, the strength of the metal lines, especially low-ionisation lines such as C II, imply that the neutral hydrogen column density is $N_{\text{H}} > 10^{18} \text{ cm}^{-2}$. Note, however, that with the exception of C IV, these lines could be contaminated by Ly α forest absorption.

Note that the total gas column density in the form of the absorbing clouds is almost independent of the interpretation of the data. Gas with a *neutral* column density of $\sim 10^{19} \text{ cm}^{-2}$ is predicted to be mostly neutral and hence to have a *total* gas column density $\sim 10^{19} \text{ cm}^{-2}$. Gas with a *neutral* column density of $\sim 10^{16} \text{ cm}^{-2}$, on the other hand, is predicted to be strongly ionised by the UV background, and hence its *total* gaseous column density will be a factor of $\sim 10^3$ greater than the neutral column density. Thus the total hydrogen column would be roughly the same as for the damped Ly α interpretation.

3 The Nature of the Gas

The three QSO sightlines may be unrepresentative of the space around this cluster. If they are representative, however, this implies that a fraction of order unity of all sightlines passing within several Mpc of this cluster would intersect a cloud of hydrogen with a neutral column density of $N_{\text{H}} \sim 10^{19} \text{ cm}^{-2}$ or greater.

In this section we hypothesise that the region surrounding the 2142–4420 cluster is optically thick in neutral hydrogen clouds. We further hypothesise that this cluster is typical of clusters at this redshift. What would be the physical consequences if our hypotheses were correct?

3.1 Gas Geometry

The absorption-line gas in cluster 2142–4420 seems to be substantially more extended spatially than the Ly α -emitting galaxies, as can be seen in Figure 1. Strong absorption is seen in QSO 2138–4427, whose sightline passes $\sim 4 \text{ Mpc}$ from the concentration of Ly α -emitting galaxies. It is also seen in QSO 2139–4433, at a redshift that places the absorption about 2 Mpc in front of the cluster. We therefore hypothesise that the cluster of Ly α sources is embedded within a much larger structure of absorbing gas. Indeed, it is suggestive that the QSO sightline passing closest to the three Ly α -emitting galaxies shows the lowest neutral hydrogen column density. This might indicate that the central region of the neutral

gas structure is hotter and more ionised than the outer regions.

If this neutral gas structure does exist, what is its geometry? With only the QSO absorption to guide us, all we can do is bracket the possibilities with three straw models:

1. Sheet model: the gas lies in a sheet, of thickness $\sim 1 \text{ Mpc}$, width $> 4 \text{ Mpc}$ and depth (along the line of sight) $\sim 4 \text{ Mpc}$. This sheet would be edge-on to our sightline.
2. Spherical model: the gas lies in a spherical halo, centred on the Ly α -emitting galaxies. The halo radius must be at least 4 Mpc .
3. Filamentary model: the cluster of galaxies lie at the intersection of a number of gas-filled filaments. Each filament is at least 4 Mpc long, and perhaps around 1 Mpc thick. Filamentary distributions are predicted by many cosmological simulations (e.g. Rauch, Haehnelt & Steinmetz 1997).

Is the region around the galaxy cluster completely full of absorbing gas? Let us assume that a fraction f of all randomly chosen sightlines passing within ~ 5 proper Mpc of the galaxy cluster would intercept at least one gas cloud with a total hydrogen cross section $> 10^{19} \text{ cm}^{-2}$. The probability of all three of our QSO sightlines showing such absorption is thus f^3 . If we require this probability to be greater than 1%, that implies $f > 22\%$.

3.2 The Mass and Density of the Neutral Gas Structure

Let us assume that the cluster really is surrounded by a ~ 5 proper Mpc-scale structure of gas with $f \sim 1$ and a typical neutral absorption-line column density of $N_{\text{H}} \sim 10^{19.5} \text{ cm}^{-2}$. What would be the consequences? The average density of absorption-line gas within the gas structure will then be $\sim 8 \times 10^{10} \text{ M}_{\odot} \text{ Mpc}^{-3}$ in proper coordinates. This is ~ 100 times greater than the typical density of neutral hydrogen at this redshift (e.g. Steidel 1990).

Given this density, the combined gas mass of all the absorption-line systems within the neutral gas structure would be $> 6 \times 10^{11} \text{ M}_{\odot}$ for the sheet and filamentary models, and $> 3 \times 10^{12} \text{ M}_{\odot}$ for the spherical model.

This overdensity of baryonic matter in the form of hydrogen clouds could be explained in two ways:

1. The volume surrounding the galaxy cluster is overdense in all forms of matter; the efficiency with which baryons form absorbing gas clouds and red galaxies is the same as elsewhere in the universe.
2. The volume is not greatly overdense, but the efficiency with which baryons form neutral gas clouds and red galaxies is enhanced in this region.

Cosmological simulations favour the second explanation. At redshift 2–3, there has not been sufficient time to assemble large matter concentrations. Many authors have used analytic methods or n -body simulations to estimate the mass of matter concentrations in the high-redshift universe. In Table 3, we have collected a sample of these

Table 3. Predicted mass two-point correlation coefficients ξ

Ref.	Model details	Co-moving radius	z	$\xi(r_0, z)$
1	SCDM: $\Omega_0 = 1.0$, $\Lambda = 0$, $h = 0.5$ (small σ_8)	$8.81 h^{-1}$ Mpc	2.3	0.06
1	SCDM: $\Omega_0 = 1.0$, $\Lambda = 0$, $h = 0.5$ (large σ_8)	$8.81 h^{-1}$ Mpc	2.3	0.17
2	$\Omega_0 = 1.0$, $\Lambda = 0.0$	11.4 Mpc	2.8	0.008
2	$\Omega_0 = 0.2$, $\Lambda = 0.0$	11.4 Mpc	2.8	0.095
2	$\Omega_0 = 0.2$, $\Lambda = 0.8$	11.4 Mpc	2.8	0.07
3	$\Omega_0 = 1.0$, $\Lambda = 0.0$	$10 h^{-1}$ Mpc	2.0	0.020
3	$\Omega_0 = 1.0$, $\Lambda = 0.0$	$10 h^{-1}$ Mpc	3.0	0.016
4	SCDM (with Zel'dovich approximation)	$10 h^{-1}$ Mpc	2.0	0.05
5	SCDM: $\Omega_0 = 1$, $\Lambda = 0$, $h = 0.5$ (small $\sigma_1 0$)	$10 h^{-1}$ Mpc	2.4	0.0036
5	SCDM: $\Omega_0 = 1$, $\Lambda = 0$, $h = 0.5$ (large $\sigma_1 0$)	$10 h^{-1}$ Mpc	2.4	0.03
5	TCDM: $\Omega_0 = 1$, $\Lambda = 0$, $h = 0.5$ (small $\sigma_1 0$)	$10 h^{-1}$ Mpc	2.4	0.0014
5	TCDM: $\Omega_0 = 1$, $\Lambda = 0$, $h = 0.5$ (small $\sigma_1 0$)	$10 h^{-1}$ Mpc	2.4	0.0069
5	OCDM: $\Omega_0 = 0.4$, $\Lambda = 0$, $h = 0.65$	$10 h^{-1}$ Mpc	2.4	0.016
5	Λ CDM: $\Omega_0 = 0.4$, $\Lambda = 0.6$, $h = 0.65$	$10 h^{-1}$ Mpc	2	0.0036
6	SCDM: $h = 0.5$	$10 h^{-1}$ Mpc	2.4	0.02
7	LCDM:	$10 h^{-1}$ Mpc	3	<0.005
7	C+HDM:	$10 h^{-1}$ Mpc	3	<0.005

(1) Brainerd & Villumsen (1994); (2) C  lin, Carlberg & Couchman (1997); (3) Matarrese et al. (1997); (4) Porciani (1997); (5) Moscardini et al. (1998); (6) Bagla (1998); (7) Ma (1999).

results, parametrised as the two-point correlation coefficients ξ for mass on the approximate scale of the gas structure, at redshifts $z \sim 2.38$.

Given these values of ξ , we can estimate the root-mean-squared (rms) fluctuations σ_R between the average densities of various spherical regions of this size in the early universe. We use the approximate relationship

$$\sigma_R^2 \approx 2.5\xi(R), \quad (1)$$

which is obtained from the relation $\sigma_R^2 = 3J_3(R)/R^3$, where a top-hat window function has been assumed and J_3 represents the integrated two-point correlation function (Kolb & Turner 1990). Thus

$$J_3 = \frac{1}{4\pi} \int_0^R \xi(r) d^3r, \quad (2)$$

where the two-point correlation function is assumed to be of the form $\xi(r) = (r/r_0)^{-\gamma}$, with $\gamma \approx 1.8$ (Groth & Peebles 1977; Davis & Peebles 1983).

Thus typical predicted rms mass fluctuations on the scale of this absorption-line structure (~ 10 co-moving Mpc) are only $\sim 25\%$. Even for models with the most extreme fluctuations (low-density models; Cen 1998), and assuming that the cluster sits in a 5σ mass fluctuation, the average density on this scale cannot be more than twice the mean density of the universe. Note that this applies to a roughly spherical volume: if the structure really is sheet-like or filamentary, the overdensity *within this structure* can be significantly greater.

So, the average mass density of the ~ 10 co-moving Mpc-scale volume including the cluster must be of the same order as that of the universe as a whole at this redshift. If, as seems likely, baryonic and non-baryonic matter trace each other on these large scales, primordial nucleosynthesis thus gives us an approximate upper

limit on the average baryon density of the cluster halo. Assuming $\Omega_{\text{baryon}} = 0.016 h_{100}^{-2}$ (e.g. Walker et al. 1991), and choosing $H_0 = 70 \text{ km s}^{-1} \text{ Mpc}^{-1}$, this density is $1.2 \times 10^{-26} \text{ kg m}^{-3}$ ($\sim 2 \times 10^{11} \text{ M}_\odot \text{ Mpc}^{-3}$).

Thus even if the QSO absorption-line column densities are representative of the whole cluster halo, the baryonic mass of the cluster inferred ($0.6\text{--}3 \times 10^{12} \text{ M}_\odot$ in the form of absorbing clouds) is physically possible: it does not exceed the predicted baryonic mass within the cluster volume ($\sim 7 \times 10^{12} \text{ M}_\odot$). The puzzle would be the high fraction of these baryons that are incorporated into absorption-line systems: the efficiency of formation of these objects is at least 10% and may well be much higher. This efficiency is far higher than is typical at this or any other redshift (e.g. Cen & Ostriker 1999).

3.3 The Physical State of the Gas

If the QSO sightlines are representative of all sightlines through the region around the cluster, it seems to be embedded in a structure of size ~ 5 Mpc, and most sightlines through this structure intersect a gas cloud with a hydrogen column density $N_{\text{H}} > 10^{19} \text{ cm}^{-2}$. Is there any physically plausible structure with these properties?

3.3.1 Argument from the Total Baryon Density

In this section, it is shown that the neutral gas structure is probably in the form of many small dense gas clouds, and a crude upper limit is placed on the size of these clouds. In summary, the argument is this: any given ~ 10 Mpc region of the early universe must have a baryon density that is close to the average for the whole universe, as discussed in Section 3.2. If these baryons were spread uniformly throughout the region, they would be highly ionised by the UV background radiation and no absorption would be seen. The baryons must therefore be confined in dense

clouds, occupying a small fraction of the total region. The density must be high enough that the recombination rate balances the rate of photoionisation by UV background photons. The region could contain a small number of large dense clouds, or a large number of small dense clouds. Only in the latter case, however, would most QSO sightlines through the region intercept one of these clouds (smaller clouds having a greater ratio of surface area to volume). We now consider this argument in detail.

As discussed in Section 3.2, the density of gas within the ~ 10 Mpc-scale volume surrounding the galaxy cluster can at most be comparable to Ω_{baryon} . If these baryons were spread uniformly throughout this volume, and exposed to the average UV background at this redshift, they would be highly ionised. Why then do the QSO sightlines show such large neutral column? The mean free path of UV photons at this redshift is ~ 500 Mpc (e.g. Haardt & Madau 1996); this is far greater than the average separation of UV sources, implying that the UV background intensity is spatially very uniform (Zuo 1992a,b; Fardal & Shull 1993). In this case, we know that at least five QSOs lie within 500 Mpc of the cluster, so if their emission is isotropic, lack of a UV background cannot explain the neutrality of the gas.

Let us therefore consider a model of the neutral gas structure which contains at most this average density of baryons. This mass, instead of being distributed uniformly, is confined in clouds of scale length r and density ρ . If the clouds are sufficiently large and dense, the gas within them will become neutral.

For any given density, a gas cloud must have a certain minimum size for the hydrogen within it to be neutral. This size was estimated analytically, using standard equilibrium photoionisation, and by using the MAPPINGS II photoionisation code (Sutherland & Dopita 1993). If a UV background with a plausible spectrum and intensity is assumed (the details make little difference to the final result), then the UV background ionises a layer of thickness r on the gas surface. The recombination rate within this surface layer, which is proportional to the square of its density ρ , must balance the photoionisation rate. Inserting numbers, we find that

$$r \sim k_0 \rho^{-2}, \quad (3)$$

where $k_0 \sim 3 \text{ pc cm}^{-6}$. Unless a cloud is thicker than this, its neutral column density will be low (e.g. Lanzetta 1991).

This gives one constraint on the size of the absorbing clouds. A second constraint comes from our assumption that the three QSO sightlines are representative of the region around the galaxy cluster. This assumption requires that a fraction f (of order unity) of all lines of sight passing close to the cluster intersect a cloud. To keep the discussion more general, let us allow for the possibility of sheet-like clouds of thickness r and face-on cross-sectional area A , where $A > r^2$, so that most of the incident UV flux enters through their face. If there are N clouds per unit volume throughout the region surrounding the galaxy cluster, and

the region has a typical thickness T , then this condition implies that

$$NAT \sim 1, \quad (4)$$

We can also constrain N , r and ρ by requiring that the total density $NrA\rho$ be at most comparable to ρ_{max} , the average baryon density of the universe at this redshift. These constraints can only be met if r is small. Solving, we find that A cancels, leaving us with the limits

$$r < \frac{\rho_{\text{max}}^2 T^2}{k_0} \sim 1 \text{ kpc} \quad (5)$$

and

$$\rho > \frac{k_0}{\rho_{\text{max}} T} \sim 0.03 \text{ cm}^{-3} \quad (6)$$

If the absorption is caused by approximately spherical clouds (i.e. $A \sim r^2$), these clouds could have masses of $\sim 10^6 M_{\odot}$ or less, and a space density of $N > 10^5 \text{ Mpc}^{-3}$. If they are flattened, they could be more massive and rarer.

3.4 Argument from Metal-line Ratios

Most high-column-density QSO absorption-line systems, when observed with sufficiently high spectral resolution, break down into multiple components with different ionisation states (e.g. Prochaska & Wolfe 1997). The wavelength shifts between different metal-line species seen in our data suggest that the same is true for our Lyman-limit systems.

Our spectra are not of sufficient resolution to resolve this substructure, so we are forced to use single-cloud modelling. Nonetheless, useful constraints can be put on the cloud size. We used the MAPPINGS II photoionisation code, as before, to estimate absorption column densities as a function of cloud density and metallicity.

No useful constraints could be placed on the metallicity of the gas, other than noting that it contains metals. The presence of strong high-ionisation lines, particularly C IV and Si III, however, implies that a large part of the cloud mass cannot be at densities significantly higher than $\sim 0.1 \text{ cm}^{-3}$. The Si III measurement could be contaminated by Ly α forest lines, but C IV should be reliable.

Given this maximum density, the UV background ionises a layer ~ 1 kpc deep into the cloud surface. This places a lower limit on the scale length r of the absorbing clouds, if the high- and low-ionisation lines come from the same clouds. This limit is only marginally consistent with the upper limit placed on cloud size in Section 3.3.1.

Alternatively, the absorption-line systems could be a mixture of these low-density clouds, and much smaller, higher-density neutral clouds that contribute most of the low-ionisation and Ly α absorption.

3.5 Argument from the Velocity Dispersions

Are the hypothesised absorbing clouds spread uniformly throughout the region around the galaxy cluster, or are

they gathered into clumps (Figure 4)? The velocity structure of the absorption lines suggests that the latter is true. The absorption (Figure 2) consists of a number of discrete components; the velocity width of each component ($b \sim 100 \text{ km s}^{-1}$) is much smaller than the velocity dispersion between the different components ($\sim 6000 \text{ km s}^{-1}$).

Each of the individual Ly α absorption components can only be well fitted if the absorbing gas has a velocity dispersion of $b \sim 100 \text{ km s}^{-1}$. (If a velocity dispersion of $b < 50 \text{ km s}^{-1}$ is used in the fits, the column density has to be very large to match the width of the base of the absorption component, in which case the predicted damping wings are much larger than observed.) The velocity dispersions of each component, while smaller than those of the cluster as a whole, are thus greater than the probable thermal velocity widths of the clouds themselves ($b < 30 \text{ km s}^{-1}$ for a hot photoionised phase).

Why are the individual absorption components so broad? As discussed in Section 3.4, each component is probably made up of many small absorbing clouds. If these small clouds are not gravitationally bound, the velocity dispersions could be Hubble-flow redshift differences between the clouds at the front and back of the clump; this would imply clump sizes of $\sim 100 \text{ kpc}$. If, however, they are gravitationally bound, we can use the virial theorem to constrain the mass of each clump, and hence put a lower limit on the size of each clump by requiring that the total mass of all the clumps in the cluster not exceed a cosmologically plausible amount.

Assume that each clump has size l and mass M_c . Given the observed velocity dispersions σ_v , the total mass of each clump can be estimated from the virial theorem, and

is proportional to the clump size:

$$M_c \sim \frac{l \sigma_v^2}{G}. \quad (7)$$

As in Section 3.3.1, we can constrain the size l and space density N_c of these clumps by requiring that most sightlines through the cluster must intersect a clump, i.e. that $f \sim 1$. Thus

$$N_c l^2 T \sim 1. \quad (8)$$

In addition, the total mass of all the clumps within a given volume, $N_c M_c$, can at most be comparable to the critical density of the universe ρ_{crit} , as discussed in Section 3.2. Thus the total mass in all the clouds must be inversely proportional to the typical cloud size l . Solving, we find that

$$l > \frac{\sigma_v^2}{G T \rho_{\text{crit}}} \sim 20 \text{ kpc}. \quad (9)$$

If the clumps were smaller, each would be less massive, but so many would be needed to intersect most QSO sightlines that the total mass of all the clumps would be physically implausible.

Thus each absorption component probably consists of many small absorbing clouds gathered into clumps of size $> 20 \text{ kpc}$.

3.6 Discussion

If we assume that most sightlines passing close to the 2142–4420 cluster intersect a cloud with $N_{\text{H}} \sim 10^{19} \text{ cm}^{-2}$, and that this cluster is typical of high-redshift clusters, these clusters must be embedded in

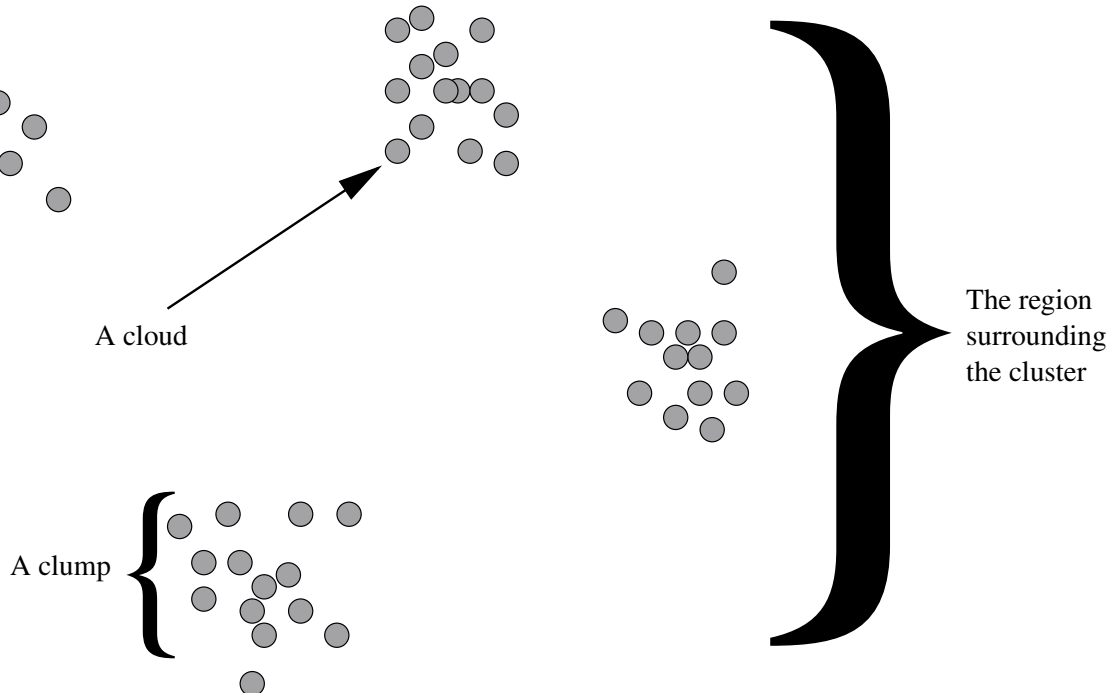


Figure 4 Diagram illustrating our terminology.

~ 10 co-moving Mpc structures of neutral gas, made up of large numbers of small dense gas clouds, gathered into clumps.

If these clumps exist, what could they be? Each clump could be the halo or extended disk of a galaxy (e.g. Prochaska & Wolfe 1997): if they are ~ 100 kpc in size, the virial mass of each halo would be $\sim 10^{11} M_{\odot}$, and the whole cluster halo would have to contain a few hundred of these galaxy halos. Alternatively, the velocity dispersions could represent infall into the potential wells of proto-galaxies (e.g. Rauch, Haehnelt & Steinmetz 1997). The properties we infer for the absorbing gas within the cluster are strikingly similar to the theoretical predictions of lumpy gas infalling along filaments into protogalaxies (e.g. Mo & Miralda-Escudé 1996).

If the halo model is correct, we should see galaxies associated with the absorption systems. For the Lyman-limit absorber in the spectrum of QSO 2139–4434, we do see a galaxy that could plausibly be surrounded by a halo that is causing the absorption: one of the strong Ly α sources lies at the same redshift as the absorber, $20''$ (~ 100 projected kpc) from the QSO sightline. For the damped systems in QSO 2139–4433 and QSO 2138–4427, however, no galaxies are seen within 500 projected kpc of the sightlines. This is not, however, conclusive, as any galaxies could easily lack Ly α emission signatures, or lie below our flux limits.

When observed at high angular resolution, many radio galaxies are seen to be surrounded by numerous kpc-sized knots of emission, sometimes made visible by reprocessing nuclear emission (e.g. Stockton, Canalizo & Ridgway 1999; Pascarella et al. 1996). These knots may be the same compact gas clouds we are hypothesising, based on the absorption.

Could all damped Ly α systems be caused by absorption in galaxy proto-clusters? The space density of Lyman-break galaxy concentrations (Steidel et al. 1998) is comparable to that of damped Ly α systems, so if all these concentrations are optically thick throughout in Ly α , then most damped Ly α systems could indeed arise in proto-clusters. Note, however, that our lower limit on the size of these clumps of small absorbing clouds is greater than the tentative upper limit placed on the size of two damped Ly α systems by Møller & Warren (1998) and Fynbo, Møller & Warren (1999).

What could be the physical origin of the absorbing clouds? Thermal instabilities can produce small cold dense clouds within the halos of high-redshift galaxies (e.g. Viegas, Friaça & Gruenwald 1999) or in collapsing proto-galaxies (Fall & Rees 1985). It is intriguing that our inferred cloud masses are comparable to those of globular clusters. The sound crossing times of these small clouds would be small, so if they are long-lived objects, they would need to be confined in some way.

Given that all sightlines though a clump intersect multiple absorbing clouds, then the mean time between cloud–cloud collisions would be less than the time it takes for clouds to cross the clump, i.e. $100 \text{ kpc}/100 \text{ km s}^{-1} \sim 10^9 \text{ yr}$

(cf. McDonald & Miralda-Escudé 1999), further evidence for the transient nature of these clouds. What happens to the gas in the clouds during a collision? The collision timescale is $\sim 1 \text{ kpc}/100 \text{ km s}^{-1} \sim 10^7 \text{ yr}$; if all the gas were converted into stars, this would imply a star formation rate of $\sim 0.1 M_{\odot} \text{ yr}^{-1}$, which is slightly below the detection threshold of most current surveys for high-redshift star-forming galaxies. Over the neutral gas structure as a whole, however, the integrated effect of all the cloud–cloud collisions, if they all form stars, is a star-formation rate of $\sim 10^3 M_{\odot} \text{ yr}^{-1}$, which should produce a diffuse H α flux of $\sim 3 \times 10^{-19} \text{ erg cm}^{-2} \text{ s}^{-1} \text{ arcsec}^{-2}$ (Kennicutt 1983) over the entirety of the cluster, and in the absence of self-absorption, a comparable flux of diffuse Ly α emission. We address the detection of such faint Ly α fluxes in Section 4.

4 Diffuse Ly α Emission

How can we observationally determine whether neutral gas structures like the ones hypothesised in Section 3 exist? One approach would be to find larger samples of high-redshift galaxy clusters with background QSOs, but this would be enormously expensive in telescope time. Another approach is to search for diffuse Ly α emission from the hypothesised neutral gas.

Even in the absence of photoionisation from young stars and/or AGN, the strong UV background in the high-redshift universe should photoionise the outer layers of any neutral gas clouds in the early universe, and hence induce Ly α emission. This emission has been modelled by Binette et al. (1993) and by Gould & Weinberg (1996) (see also Hogan & Weymann 1987). Predicted Ly α surface brightnesses are very low, typically $10^{-19} \text{ erg cm}^{-2} \text{ s}^{-1} \text{ arcsec}^{-2}$ or less. Nonetheless, if the emission covers a large area on the sky, the sensitivity can be increased by \sqrt{a} , where a is the area integrated over. This approach has been applied to emission from individual Lyman-limit systems, and stringent upper limits placed on the surface brightness (Bunker, Marleau & Graham 1998).

We used results from these papers, together with our own modelling using MAPPINGS II, to predict the diffuse Ly α flux from the hypothesised neutral gas structure around the 2142–4420 cluster. The biggest sources of uncertainty are the strength of the UV ionising background at this redshift (literature values disagree by a factor of ~ 3), the spectrum of the UV background, and the geometry of the structure (which can also introduce a factor of ~ 3 uncertainty into the predicted diffuse flux). All predictions ignore dust obscuration; if any dust is present, then the Ly α fluxes will probably be reduced by a large factor. We predicted that the diffuse flux per unit area F should lie in the range $10^{-17} > F > 5 \times 10^{-20} \text{ erg cm}^{-2} \text{ s}^{-1} \text{ arcsec}^{-2}$.

4.1 Observations

We have made two attempts to detect diffuse Ly α emission from the hypothesised neutral gas structure surrounding the 2142–4420 cluster.

The first attempt used the Double Beam Spectrograph on the Siding Spring 2.3 m telescope. A long-slit spectrum was taken, with the slit centred on QSO 2139–4434, and the slit angle positioned so that it included QSO 2139–4433. The slit was 2'' wide and 3' long, giving a total sky coverage of 360 square arcsec. Spectral resolution (2 pixels) was 1.1 Å. 1800 s exposures of the cluster field were alternated with equal length exposures on randomly chosen control fields 1° away (in random directions).

A total on-field exposure time of 23,400 s was obtained, in dark photometric conditions, on the nights of 1999 July 9 and 10. The total control field exposures were identical. The night sky spectrum was found to be very stable throughout both nights, so on- and off-field exposures were combined, using inverse variance weighting to minimise sky noise. The night sky brightness at around 4110 Å (the expected wavelength of Ly α at the cluster redshift) was $\sim 7 \times 10^{-18}$ erg cm $^{-2}$ Å $^{-1}$ arcsec $^{-2}$.

No significant excess flux was seen at wavelengths of around 411 nm in the on-field image as compared to the off-field image. This enables us to place a 3 σ upper limit on the diffuse emission from the cluster halo, if it fills the slit, of 1.8×10^{-18} erg cm $^{-2}$ s $^{-1}$ arcsec $^{-2}$.

The second attempted measurement was made with the Taurus Tunable Filter (TTF), on the Anglo-Australian Telescope. TTF is an imaging Fabry-Perot etalon system, with a resolution at 4100 Å of down to 4 Å, and a monochromatic field of view approximately 5' in diameter. All observations were taken in photometric conditions on the nights of 1999 August 11 and 12.

Charge shuffling was used to alternate, every 60 s, between images of the cluster field at the expected wavelength of Ly α , and at a control wavelength 50 Å to the red. Observations were made at 4 Å resolution, with central wavelengths of 4110 Å and 4114 Å. On-wavelength exposure times were 3000 and 3600 s respectively. The field of view was centred on the brightest of the Ly α emitting galaxies, at 21:42:27.48–44:20:28.4 (J2000).

The difference images (on-wavelength minus off-wavelength) were very clean, with all continuum sources removed to high precision. These images were aligned and co-added, using inverse variance weighting. As no control fields were observed, we have no sensitivity to diffuse flux filling the field of view. We are sensitive to variations in the diffuse flux on smaller angular scales. On scales of 20'' \times 20'', we see no regions of excess flux, to a 3 σ upper limit of 5×10^{-19} erg cm $^{-2}$ s $^{-1}$ arcsec $^{-2}$, at both wavelengths.

4.2 Summary

We failed to detect diffuse emission from any neutral gas around the 2142–4420 cluster. Our observations were not, however, deep enough to rule out the presence of even dust-free neutral gas. They do, however, demonstrate the darkness and stability of the sky background at these blue wavelengths.

5 Conclusions

The QSO absorption-line data on the 2142–4420 cluster are provocative. If taken at face value, they suggest that this high-redshift cluster, and perhaps all of them, are surrounded by ~ 5 Mpc structures of neutral hydrogen, of total mass $\sim 10^{12}$ M $_{\odot}$. This gas would be gathered into a series of >20 kpc clumps, each of which is made up of hundreds of small (<500 pc), dense (>0.1 cm $^{-3}$) clouds. These clouds may merge to form cluster galaxies, collapse to form globular clusters, or perhaps dissolve into X-ray intra-cluster medium by redshift zero.

Many more observations of neutral gas around high-redshift clusters will be needed to verify this picture.

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