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Influence of A-site Cation on Structure and Dielectric Properties in A_2 DyBiO₆ (A = Mg, Ca, Sr, Ba) Double Perovskites

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Double perovskite metal oxides with formula A_2DyBiO_6 (A=Mg, Ca, Sr, Ba) were synthesised by a sol–gel autocombustion method, using citric acid as the combustion agent. The effects of A-site cation on the structure, morphology, and dielectric properties were examined. The synthesis was monitored using Fourier transform infrared spectroscopy (FTIR) to indicate the absence of organic phase. X-ray diffraction (XRD) analysis showed that the compounds have three different perovskite structures. Structural characterisation of the samples was evaluated using XRD patterns. Scanning electron microscopy showed that all samples are formed by agglomerated particles. Dielectric properties were evaluated using dielectric permittivity and dielectric losses. Cole–Cole plots show a single semicircle for all materials, indicating that the double perovskites obtained are composed of well conducting grain boundaries and poorly conducting grains.

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Introduction

Double perovskites, with $A_2BB'O_6$ formula, are formed by doubling the unit cell of a simple perovskite-type, ABO_3 . These materials have been intensively studied in recent years, especially due to their capacity to incorporate many elements such as alkaline earth (A-site), transition metal, or lanthanide cations (B, B'-site) in their structures. Therefore, the structure of these materials can change due to octahedral tilting (BO₆, B'O₆) to build a double perovskite-type lattice. [1] This leads to the possibility of preparing materials for specific applications. These features are very important for many magnetic materials [2,3] and electronic devices that can be operated successfully in a wide range of frequencies. [4] These double perovskites are also used in various fields such as catalysis, [5] photocatalysis, [6] and solid oxide fuel cell (SOFC) materials.

Few reports have been made about double perovskites with Bi-cation on the B^\prime -site. Byeon et al. $^{[8]}$ studied the synthesis and characterisation of $Ba_2Bi^VFe^{III}O_6$. Macquart and Kennedy $^{[9]}$ investigated the structure of $Ba_{2-x}Sr_xLuBiO_6$ double perovskite. Harrison et al. $^{[10]}$ reported the synthesis, structural characterisation, and magnetic properties of Ba_2MBiO_6 (M = Ce to Lu) double perovskites. Lenz and Müller-Buschbaum $^{[11]}$ studied the structures of Ba_2BiMO_6 , (M = Y, Dy) double perovskites and concluded that these compounds exhibit a cubic structure with a Fm-3m space group. Horyń et al. $^{[12]}$ studied Sr_2REBiO_6 (RE = La and lanthanides) double perovskites and reported that Sr_2DyBiO_6 has a monoclinic structure with a $P2_1/n$ space group.

In all of these studies, the compounds were prepared using a solid-state reaction method, except for the synthesis of Ba₂Bi^VFe^{III}O₆ by Byeon et al.,^[8] which was carried out under a high pressure oxygen atmosphere.

In a previous paper, we studied the compound, Ca_2DyBiO_6 . ^[13] Electrical measurements revealed that this compound was the best dielectric material, displaying a dielectric constant of 40 with very low dielectric loss. Therefore, starting from Ca_2DyBiO_6 , the A-site cation was substituted in order to improve its electrical properties.

Hence, the aim of the present work is to study, for the first time, the effect of the A-site cation on the structure and dielectric properties in A_2DyBiO_6 (A = Mg, Ca, Sr, Ba) double perovskite-type oxides. An important technique for double perovskite preparation is the sol–gel method because this procedure allows the control of chemical composition, solid texture, size, and shape of the particles. [14] In this study, the materials are obtained by a sol–gel auto-combustion method using citric acid as a chelating and fuel agent. Our results showed that the dielectric properties were improved with larger ionic radius cation substitution at the A site.

Experimental

Nanopowders of A_2DyBiO_6 (A=Mg, Ca, Sr, Ba) double perovskite-type oxides were prepared in air using a sol-gel auto-combustion method with citric acid $C_6H_8O_7$ (Merck)

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as a chelating and fuel agent. Analytical grade precursors, $Mg(NO_3)_2 \cdot 4H_2O$ (Merck), $Ca(NO_3)_2 \cdot 4H_2O$ (Merck), $Sr(NO_3)_2 \cdot 4H_2O$ (Merck), $Ba(NO_3)_2 \cdot 4H_2O$ (Merck), $Dy(NO_3)_3 \cdot 6H_2O$ (Merck), and $Bi(NH_3)_2C_6H_7O_7 \cdot H_2O$ (Merck) were mixed in stoichiometric proportions. A solution of citric acid was mixed with metal nitrate mixtures in a 3:1 molar ratio of citric acid to metallic cation. The solutions obtained were brought to an alkaline pH to avoid precipitation. For A_2DyBiO_6 (A=Mg, Ca, Sr, Ba) double perovskite-type oxides, the following chemical equation could be written:

$$\begin{split} 2A(NO_3)_2 + Dy(NO_3)_3 + Bi(NH_3)_2C_6H_7O_7 \cdot H_2O \\ + (37/2)O_2 + 3C_6H_8O_7 &= A_2DyBiO_6 + 9NO_2 \uparrow \\ + 12CO_2 \uparrow + (39/2)H_2O \uparrow \end{split} \tag{1}$$

Eqn 1 is a proposed equation following a stoichiometric calculation that takes into account the redox process. The required amounts of metal salts to obtain the desired compounds have not been determined based on Eqn 1, but based on the mass reports, considering the molecular mass of general formula proposed.

The mixed solutions were heated up to 353 K under continuous stirring over a water bath in order to obtain viscous gels. The dried gels were gradually heated on a sand bath at 623 K for 1 h, when the auto-combustion was clearly observed and all the gels were converted to powders. The obtained powders were thermally treated in three steps: up to 773 K for 7 h, up to 973 K for 7 h, and up to 1173 K for 9 h.

Infrared (IR) and X-ray diffraction (XRD) analyses were recorded in powder form at different sintering temperatures. The progress of the synthesis and the disappearance of the organic phase were monitored by IR spectroscopy using a Bruker spectrophotometer TENSOR TM 27-type with Fourier transform with an ATR cell at 2 cm $^{-1}$ resolution in the IR range of 4000 to $400\,\mathrm{cm}^{-1}$. The double perovskite phase formation was established using powder X-ray diffraction, recorded by means of a Shimadzu LabX 6000 diffractometer equipped with graphite monochromator and CuKa ($\lambda = 1.5406\,\mathrm{\mathring{A}}$) radiation. The samples fixed in the reflection mode were analysed under ambient atmosphere with a scanning rate of $0.02^{\circ}\,\mathrm{s}^{-1}$ over the $2\theta = 20-80^{\circ}$ range.

The morphology of the sintered powders at 1173 K was investigated by scanning electron microscopy (SEM) using a Vega Tescan-type scanning transmission electron microscope. The micrographs were made in fracture, i.e. powders heated at 1173 K were pressed into pills without other thermal treatment.

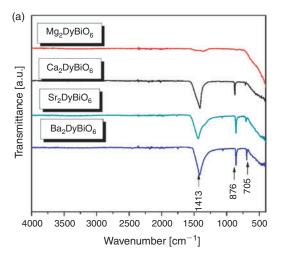
The dielectric measurements were carried out using the twoterminal direct current method. The dependence of dielectric permittivity and dielectric losses on frequency was studied using an Agilent 4292–1 device, in the range of 40 Hz to 10 MHz. The powders were thermally treated at 1173 K, pressed into disks at 400 kPa cm⁻² without subsequent calcinations, and were inserted between two electrodes.

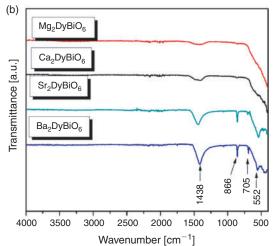
Results and Discussion

IR Results

IR spectra for the A₂DyBiO₆ samples treated at 773, 973, and 1173 K, recorded at room temperature, are presented in Fig. 1a–c.

IR spectra of pure metal oxide samples are very similar, showing two bands characteristic to the symmetric (ν_a) and antisymmetric (ν_a) octahedral BiO₆ stretching vibrations. The shift of stretching vibration to lower wavelengths is due to stronger





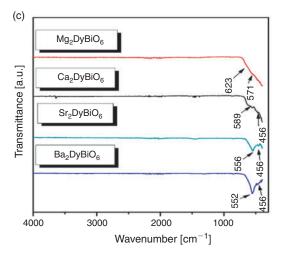


Fig. 1. IR spectra of $\rm A_2DyBiO_6$ samples treated at (a) 773 K, (b) 973 K, and (c) 1173 K.

interactions between the cation and oxygen atoms (owing to the electronegativity of the metal). $^{[15,16]}$

Therefore, the IR spectra at 773 (Fig. 1a) and 973 K (Fig. 1b) depict the specific symmetric (ν_s) and anti-symmetric (ν_{as}) bands for metal–oxygen stretching vibration in the range of 900 to $400\,\mathrm{cm}^{-1}$ and the absorption peak of nitrate groups at around $1417\,\mathrm{cm}^{-1}$.

Furthermore, the IR spectra for A₂DyBiO₆ samples heated at 1173 K (Fig. 1c) showed the disappearance of the nitrate groups.

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Note the decrease in intensity of the symmetric (υ_s) band with the introduction of cations with smaller ionic radii in the system. Also, the anti-symmetric υ_{as} stretching vibration (589 to 552 cm $^{-1}$) of double perovskite-type oxide^[15] characteristic of octahedral BiO₆, and the bands at 437 cm $^{-1}$ characteristic of Dy–O stretching vibration,^[17] become predominant with increasing temperature. This asymmetric band intensity variation leads to the idea that the Bi–O bond strength increases with increasing ionic radii of A²⁺cations, which means stabilisation of compound structures.

The presence of two lower bands at 623 and 571 cm⁻¹ can also be seen in the IR spectra of Mg₂DyBiO₆, which can be attributed to the characteristic bands of a pyrochlore structure.^[18]

XRD Results

XRD patterns of A_2DyBiO_6 materials heated at $1173\,K$ are presented in Fig. 2a–d. According to literature data, Ca_2DyBiO_6 and Sr_2DyBiO_6 double perovskites were found to have monoclinic structures with $P2_1/n^{[19]}$ and $P12_1/n1^{[12]}$ space groups, respectively. The strongest reflection peak at a 2θ value of around 30° identified in Ca_2DyBiO_6 and Sr_2DyBiO_6 samples are assigned to $(1\,1\,2)$. Mg_2DyBiO_6 double perovskite is indexed as a defective cubic pyrochlore structure with a Fd-3m space group, while Ba_2DyBiO_6 double perovskite crystallised in cubic system has a Fm-3m space group. Unfortunately, the existence of a secondary phase can be observed for the Mg and Ba samples and this was identified as Bi_2O_3 (cubic system, COD1010313), according to the diffraction peak at $2\theta = 27.3^\circ$.

Crystallite size (*D*) was calculated from XRD using the Debye–Scherrer^[21] equation (Eqn 2):

$$D = \frac{0.94 \cdot \lambda}{\beta_{1/2} \cdot \cos \theta} \tag{2}$$

where D = crystallite size, $\lambda =$ wavelength of X-ray source (CuK $\alpha = 1.5405$ Å), $\beta_{1/2}$ (rad) = half maximum of diffraction plane (h k l), $\theta =$ angle of diffraction.

The pseudo-cubic pyrochlore structure of the Mg_2DyBiO_6 compound contains large holes that can lead to structural defects and therefore defaults to higher crystallites. A distorted structure leads to crystallite (D) with larger size, as can be seen in Table 1.

The best ordered structure was obtained in the case of the Ba₂DyBiO₆ compound, which has smaller sizes of crystallites (28 nm).

SEM Results

Fig. 3a–d shows the SEM images for all samples heated at 1173 K. All micrographs were made in fracture (powders thermally treated at 1173 K were pressed into pills without other additional thermal treatment). It can be observed that the materials are formed by agglomerated particles. Generally, the decrease in particle size is directly proportional to the increase in specific surface of these oxides. The irregular variation of $S_{\rm BET}$ values for the studied samples (3 m² g¹ for Mg²DyBiO6, 1 m² g¹ for Ca²DyBiO6, 12 m² g¹ for Sr²DyBiO6, and 6 m² g¹ for Ba²DyBiO6) can be due to the different degrees of agglomeration of grains and number of fine pores. From the SEM image of the sample with Mg, it can be seen that the particles are crowded close together and almost 'glued' to each other. The presence of holes is most likely caused by the exhaust gases from the combustion process.

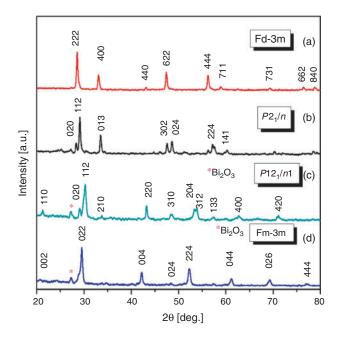


Fig. 2. XRD spectra of (a) Mg_2DyBiO_6 , (b) Ca_2DyBiO_6 , (c) Sr_2DyBiO_6 , and (d) Ba_2DyBiO_6 samples heated at 1173 K.

Table 1. Crystallite size (D) from XRD

Compound	D [nm]
Mg ₂ DyBiO ₆	48
Ca ₂ DyBiO ₆	40
Sr ₂ DyBiO ₆	32
Ba ₂ DyBiO ₆	28

Dielectric Results

Fig. 4 a, b depicts the variation of real parts (ε') and imaginary parts (ε'') of dielectric constant with frequency. The dielectric constant dependence of the frequency for all A2DyBiO6 (A = Mg, Ca, Sr, Ba) double perovskites were studied in the frequency range of 40 Hz to 10 MHz at room temperature. It can be clearly observed that all obtained materials exhibit a normal dielectric dispersion because the dielectric constants (ε' and ε'') decrease with increasing frequency. One can observe that the dielectric constant decreases with increasing frequency, reaching a constant value at high frequencies. This variation can be observed in all of the sample cases. This phenomenon is a normal dielectric dispersion caused by electron displacement as a response to an applied field or defects and oxygen vacancies, known as Maxwell–Wagner-type interfacial polarisation. [22] By changing A in A₂DyBiO₆ compounds, the dielectric constant also changes. It is observed that Ba₂DyBiO₆ exhibits higher values of dielectric constant at all frequencies. Compared with the rest of the samples, Ba₂DyBiO₆ shows a much higher value of dielectric constant and the decrease in dielectric constant as a function of frequency is much sharper. It can be seen that in the case of Mg₂DyBiO₆, the dielectric constant reached the linear plateau at low frequency, probably caused by the defects presented in the pyrochlore structure. According to the experimental range of frequencies reported in the literature, the dielectric constant (ϵ') of these defective pyrochlore materials is predominantly determined by the contribution of ionic New Dielectric Compounds 253

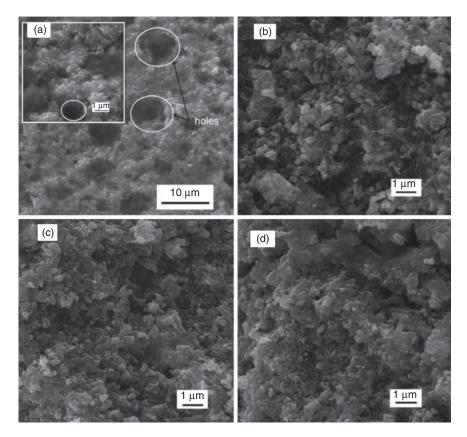


Fig. 3. SEM images for (a) Mg₂DyBiO₆, (b) Ca₂DyBiO₆, (c) Sr₂DyBiO₆, and (d) Ba₂DyBiO₆.

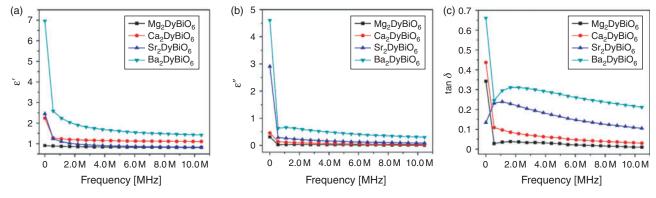


Fig. 4. Variation of (a) real part of permittivity, (b) imaginary part of permittivity, and (c) dielectric loss tangent with frequency.

polarisation. ^[20] In other double perovskite structure-containing samples, the dipole follows the field at low frequencies, but as the frequency increases, the dipoles start lagging behind the field and the real part (ϵ') of the dielectric constant decreases. At high frequencies, the dielectric constant becomes independent of frequency, as reported by Laishram et al., ^[23] because the dipoles can no longer follow the field.

The experimental data indicate that by changing A from Ca to Sr and Ba the permittivity increases, which may be related to the decrease in concentration of dipoles^[20] due to the increase in unit-cell volume.

The variation of dielectric loss values with log (frequency) of all materials are shown in Fig. 4c. It can be seen that the dielectric loss values range from 0.03 to 0.68 at $10\,\mathrm{MHz}$. It is clear that $\mathrm{Mg_2DyBiO_6}$ has smaller loss values throughout the range of frequencies. In perovskites, an important contribution

to high losses is attributed to oxygen vacancies, usually present in such systems. The increase in the real part of dielectric constant values and of $\tan \delta$, particularly at lower frequencies, by changing A from Ca to Sr and Ba, indicates important contributions from conductivity relaxation combined with Maxwell–Wagner phenomena [23] to the total dielectric response.

All double perovskite samples show values of permittivity ranging between one and seven, and low losses at room temperature in the whole range of frequencies, indicating good dielectric character of these materials compared with that from an earlier report by Vijayakumar et al.^[24]

Fig. 5a–d shows the Cole–Cole plots for all compounds obtained by plotting the variation of imaginary part (ϵ'') versus real part (ϵ') of the dielectric constant. It can be observed that the Cole–Cole plots for all samples are characterised by one semicircle. It can be noted that the grain boundary resistance

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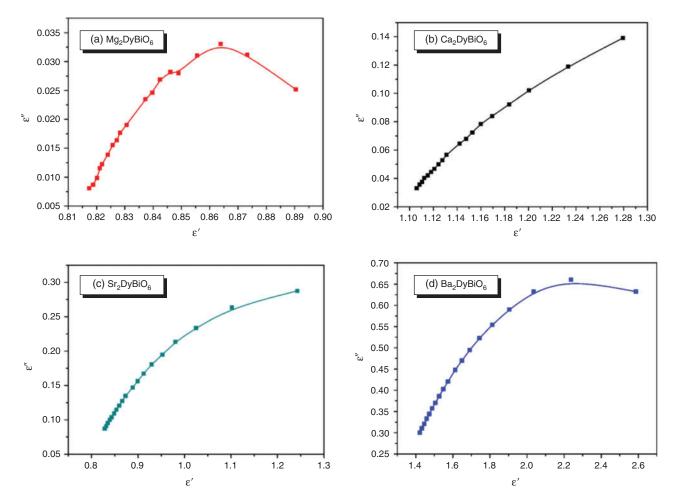


Fig. 5. Cole–Cole plots for (a) Mg_2DyBiO_6 , (b) Ca_2DyBiO_6 , (c) Sr_2DyBiO_6 , and (d) Ba_2DyBiO_6 obtained by plotting the variation of the imaginary part (ϵ'') versus the real part (ϵ') of the dielectric constant.

contributed to the dielectric measurements and is directly proportional to the polarisability of the studied samples. From the variation of the dielectric constant, it can be observed that the values of grain boundary resistance increase with increasing divalent cation (A^{2+}) ionic radius.

Conclusions

 A_2DyBiO_6 (A = Mg, Ca, Sr, Ba) double perovskites were successfully prepared by a sol–gel auto-combustion method using citric acid as a combustion agent.

The IR spectra showed bands characteristic of metal—oxygen bonds in the perovskites. XRD revealed the formation of double perovskite phases at 900°C containing the minor impurity Bi₂O₃ in Sr₂DyBiO₆ and Ba₂DyBiO₆ double perovskites. Ca₂DyBiO₆ and Sr₂DyBiO₆ were found to have monoclinic structures, whilst Ba₂DyBiO₆ had a cubic structure, and Mg₂DyBiO₆ was found to have a defective pseudo-cubic pyrochlore structure. SEM images revealed the formation of holes in the structure of Mg₂DyBiO₆ double perovskite, probably caused by the gases from the combustion process. Microstructures of Ca₂DyBiO₆, Sr₂DyBiO₆, and Ba₂DyBiO₆ compounds showed homogenous crowded particles.

All double perovskite samples showed values of permittivity ranging between one and seven and low losses at room temperature in the whole range of frequencies, indicating good dielectric character of these materials. Cole—Cole plots show a single semicircle for all materials, which means that all polycrystalline

double perovskites are composed of well-conducting grain boundaries and poorly conducting grains.

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References

- R. Shaheen, J. Bashir, Solid State Sci. 2010, 12, 1496. doi:10.1016/ J.SOLIDSTATESCIENCES.2010.06.015
- [2] C. M. Bonilla, D. A. L. Téllez, J. A. Rodríguez, E. V. López, J. Roa-Rojas, *Physica B* 2007, 398, 208. doi:10.1016/J.PHYSB.2007.04.078
- B. García-Landa, C. Ritter, M. R. Ibarra, J. Blasco, P. A. Algarabel,
 R. Mahendiran, J. García, Solid State Commun. 1999, 110, 435.
 doi:10.1016/S0038-1098(99)00079-4
- [4] R. V. K. Mangalam, E. Suard, A. Sundaresan, *Physica B* 2009, 404, 154. doi:10.1016/J.PHYSB.2008.10.031
- [5] R. Hu, R. Ding, J. Chen, J. Hu, Y. Zhang, Catal. Commun. 2012, 21, 38. doi:10.1016/J.CATCOM.2012.01.008
- [6] T. Hatakeyama, S. Takeda, F. Ishikawa, A. Ohmura, A. Nakayama, Y. Yamada, A. Matsushita, J. Yea, J. Ceram. Soc. Jpn. 2010, 118, 91. doi:10.2109/JCERSJ2.118.91
- [7] D. Marrero-López, J. Peña-Martínez, J. C. Ruiz-Morales, M. Gabás, P. Núñez, M. A. G. Aranda, J. R. Ramos-Barrado, *Solid State Ionics* 2010, 180, 1672. doi:10.1016/J.SSI.2009.11.005

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- [8] S. H. Byeon, C. G. Demazeau, H. J. Choy, C. L. Fournes, *Mater. Lett.* 1991, 12, 163. doi:10.1016/0167-577X(91)90166-4
- [9] R. B. Macquart, B. J. Kennedy, *Chem. Mater.* 2005, 17, 1905. doi:10.1021/CM0486924
- [10] T. A. W. Harrison, P. J. K. Reis, J. A. Jacobson, F. L. Schneemeyer, V. J. Waszczak, *Chem. Mater.* **1995**, 7, 2161. doi:10.1021/ CM00059A025
- [11] A. Lenz, H. K. Müller-Buschbaum, J. Less-Common Met. 1990, 161,
 L15. doi:10.1016/0022-5088(90)90331-D
- [12] R. Horyń, M. Wotcyrz, A. Wojakowski, A. J. Zaleski, J. Alloy. Compd. 1996, 242, 35. doi:10.1016/0925-8388(96)02359-6
- [13] S. Feraru, P. M. Samoila, V. Nica, A. R. Iordan, M. N. Palamaru, Chem. Pap. – Chem. Zvesti 2013, 67, 1311. doi:10.2478/S11696-013-0394-9
- [14] S. Z. Tian, J. C. Zhao, C. D. Qiao, X. L. Ji, B. Z. Jiang, Mater. Lett. 2006, 60, 2747. doi:10.1016/J.MATLET.2006.01.083
- [15] A. E. Lavat, J. E. Baran, Vib. Spectrosc. 2003, 32, 167. doi:10.1016/ S0924-2031(03)00059-6
- [16] A. E. Lavat, J. E. Baran, J. Alloy. Compd. 2008, 460, 152. doi:10.1016/ J.JALLCOM.2007.06.003

[17] C. Vijayakumar, H. P. Kumar, S. Solomon, J. K. Thomas, P. R. S. Warian, A. John, *J. Alloy. Compd.* **2009a**, 480, 167. doi:10.1016/J. JALLCOM.2009.01.123

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- [18] M. Maczka, A. V. Knyazev, Yu. N. Kuznetsova, M. Ptak, L. Macalik, J. Raman Spectrosc. 2011, 42, 529. doi:10.1002/JRS.2735
- [19] A. E. Lavat, J. E. Baran, Mater. Res. 2011, 14, 472. doi:10.1590/ S1516-14392011005000066
- [20] S. Senthil Kumar, K.V.O. Nair, J. James, J. Solid State Chem. 2004, 177, 3873. doi:10.1016/J.JSSC.2004.07.033
- [21] C. A. Triana, D. A. Landiz Téllez, J. Arbey Rodríguez, F. Fajardo, J. Roa-Rojas, *Mater. Lett.* **2012**, *82*, 116. doi:10.1016/J.MATLET. 2012.05.041
- [22] P. K. Jana, S. Sudipta, B. K. Chaudhuri, J. Phys. D: Appl. Phys. 2007, 40, 556. doi:10.1088/0022-3727/40/2/033
- [23] R. Laishram, S. Phanjoubam, H. N. K. Sarma, C. Prakash, J. Phys. 1999, D32, 2151.
- [24] C. Vijayakumar, H. Padma Kumar, V. T. Kavitha, S. Solomon, J. K. Thomas, P. R. S. Wariar, J. Koshy, *J. Alloy. Compd.* **2009b**, 475, 778. doi:10.1016/J.JALLCOM.2008.08.013