# ${ }^{50} \mathrm{Cr}(\mathrm{p}, \gamma){ }^{51} \mathrm{Mn}$ Reaction from <br> $E_{\mathrm{p}}=1.45$ to 2.07 MeV 

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## Abstract

Proton resonances in the ${ }^{50} \mathrm{Cr}(\mathrm{p}, \gamma)^{51} \mathrm{Mn}$ reaction have been studied in the range $E_{\mathrm{p}}=1450-2070 \mathrm{keV}$. Gamma-ray spectra, measured at $55^{\circ}$ to the beam direction at each of the resonances at $E_{\mathrm{p}}=1451$, $1546,1580,1600,1689,1798,1830,2031,2042$ and 2067 keV with a high resolution $35 \mathrm{~cm}^{3} \mathrm{Ge}(\mathrm{Li})$ detector, have been used to derive the decay schemes of the resonant and bound levels in ${ }^{51} \mathrm{Mn}$. Decays to the bound levels at $4 \cdot 052,4 \cdot 729,4 \cdot 925,5 \cdot 073,5 \cdot 129,5 \cdot 174,5 \cdot 188,5 \cdot 203$ and $5 \cdot 506 \mathrm{MeV}$ have been observed in the proton capture reaction for the first time, while new results have been obtained for the decay of some of these levels. The reported levels at $2 \cdot 715,2 \cdot 892,3 \cdot 046,3 \cdot 051$, $4 \cdot 203,4 \cdot 887$ and $5 \cdot 223 \mathrm{MeV}$ are confirmed. The energies of the excited levels in ${ }^{51} \mathrm{Mn}$ have been obtained with an accuracy of $\pm 3 \mathrm{keV}$ and the $Q$ value for the reaction ${ }^{50} \mathrm{Cr}(\mathrm{p}, \gamma)^{51} \mathrm{Mn}$ has been found to be $5 \cdot 269 \pm 0 \cdot 002 \mathrm{MeV}$. From angular distribution measurements for the resonances at $E_{\mathrm{p}}=2042$ and 2067 keV , spins of $5 / 2$ are proposed for each of the excited states at $7 \cdot 270$ and 7. 295 MeV .

## Introduction

The properties of the excited levels in ${ }^{51} \mathrm{Mn}$ were first investigated by Arnell (1962) from the reaction ${ }^{50} \mathrm{Cr}(\mathrm{p}, \gamma){ }^{51} \mathrm{Mn}$. He was able to identify 9 resonances in the energy range $E_{\mathrm{p}}=900-1220 \mathrm{keV}$, using NaI detectors. Later work by Arnell and Sterner (1964) at 6 resonances for the same reaction in a similar range of $E_{\mathrm{p}}$ (and again using NaI detectors) established the existence of 11 levels below $4 \cdot 0 \mathrm{MeV}$ in ${ }^{51} \mathrm{Mn}$ and provided the first information on the decay properties of these levels. Subsequently, energy values to a better accuracy for 11 levels below 3.0 MeV excitation in ${ }^{51} \mathrm{Mn}$ were reported by Sterner et al. (1966), using $0.5 \mathrm{~cm}^{3} \mathrm{Ge}(\mathrm{Li})$ and NaI detectors. Wall and Erlandsson (1967) investigated the $\gamma$-ray decay schemes of 5 resonances for the ${ }^{50} \mathrm{Cr}(\mathrm{p}, \gamma)^{51} \mathrm{Mn}$ reaction in the energy region $E_{\mathrm{p}}=1050-1970$ keV . Excited levels up to $6 \cdot 052 \mathrm{MeV}$ in ${ }^{51} \mathrm{Mn}$ with an energy resolution of 30 keV were reported by Rapaport et al. (1967) from a study of the ${ }^{50} \mathrm{Cr}\left({ }^{3} \mathrm{He}, \mathrm{d}\right){ }^{51} \mathrm{Mn}$ reaction at 12.0 MeV bombarding energy.

The principle source of information on the level decay scheme of ${ }^{51} \mathrm{Mn}$ has been the work of Forsblom et al. (1972). Those authors used a $\mathrm{Ge}(\mathrm{Li})$ detector to study 19 resonances from the ${ }^{50} \mathrm{Cr}(\mathrm{p}, \gamma){ }^{51} \mathrm{Mn}$ reaction in the energy range $E_{\mathrm{p}}=1000-1700$ keV and reported levels up to $5 \cdot 585 \mathrm{MeV}$ excitation in ${ }^{51} \mathrm{Mn}$. Spins were assigned for some of the levels by Erlandsson (1967) and Forsblom et al. (1972) from measurements of the angular distributions. In addition, $l$-value assignments were made for
certain levels from studies of the ${ }^{50} \mathrm{Cr}\left({ }^{3} \mathrm{He}, \mathrm{d}\right)^{51} \mathrm{Mn}$ reaction by Rapaport et al. (1967) and the reaction ${ }^{50} \mathrm{Cr}(\mathrm{d}, \mathrm{n}){ }^{51} \mathrm{Mn}$ by Nilsson and Erlandsson (1970).

As part of a program to investigate the decay schemes of radiative capture resonances with large $\mathrm{Ge}(\mathrm{Li})$ detectors, the present report deals with a study of the ${ }^{50} \mathrm{Cr}(\mathrm{p}, \gamma){ }^{51} \mathrm{Mn}$ reaction in the energy range $E_{\mathrm{p}}=1450-2070 \mathrm{keV}$. The aims of this work were to look for any new resonances and bound levels, to verify the existence of certain doubtful levels, to derive the decay properties of the bound levels and to locate and study isobaric analogue states.

## Experimental Procedure

The research work was carried out at two laboratories: The 2 MeV Van de Graaff accelerator of the Australian National University was used to locate the resonances and measure the $\gamma$-ray spectra, while the angular distributions were measured with the 4 MeV Van de Graff accelerator at the Centre de Recherches Nucleaires, Strasbourg.

Enriched ${ }^{50} \mathrm{Cr}$ targets were used. To minimize the contamination of occluded gases (especially ${ }^{19} \mathrm{~F}$ ), blanks of tungsten were heated in vacuum with an induction heater before being used as targets backings. Chromium powder $92 \%$ enriched with ${ }^{50} \mathrm{Cr}$ (obtained from the Oak Ridge National Laboratory) was evaporated onto the tungsten backings to produce targets of thicknesses 10 and $20 \mu \mathrm{gcm}^{-2}$. The thinner target was used for the determinations of the yield curves for the ${ }^{50} \mathrm{Cr}(\mathrm{p}, \gamma){ }^{51} \mathrm{Mn}$ reaction, while the thicker target was used to measure $\gamma$-ray spectra at different angles to obtain the angular distributions. The targets were water-cooled and a liquid nitrogen trap was placed in front of the target chamber to remove condensable vapours and to minimize carbon buildup.

A $35 \mathrm{~cm}^{3} \mathrm{Ge}(\mathrm{Li})$ detector with a resolution of better than 3 keV for the 1173 keV $\gamma$-ray from ${ }^{60} \mathrm{Co}$ was placed at $55^{\circ}$ to the beam direction to locate the resonances and to accumulate $\gamma$-ray spectra at each resonance. The spectra were stored in a 4096channel Nuclear Data ND 2200 analyser. The spectra measured at $55^{\circ}$ yielded the decay schemes and the branching ratios of both the resonances and the bound levels. The intensities of the $\gamma$-rays were derived from the calibrated intensity of the $\mathrm{Ge}(\mathrm{Li})$ detector using the 992 keV resonance from the ${ }^{27} \mathrm{Al}(\mathrm{p}, \gamma)^{28} \mathrm{Si}$ reaction and the $\gamma$-rays from a ${ }^{56} \mathrm{Co}$ radioactive source. The relative intensity errors vary from $5 \%$ for the prominent peaks to $15 \%$ for the smaller peaks. Where the decay scheme of a bound level could be determined from more than one resonance, the transition ratios were found to be consistent within experimental error.

The angular distributions were measured with an $80 \mathrm{~cm}^{3} \mathrm{Ge}(\mathrm{Li})$ detector at five angles, $0^{\circ}, 30^{\circ}, 45^{\circ}, 60^{\circ}$ and $90^{\circ}$. A fixed smaller volume $\mathrm{Ge}(\mathrm{Li})$ detector was used to normalize the angular distribution data. The results have been analysed in terms of even-order Legendre polynomials up to fourth order. For each spin combination, a value was computed of the sum of the squared and weighted differences between the experimental and theoretical number of counts divided by the number of free parameters. A computer program performed in the first step the normalization of the theoretical distributions and determined in the second step the best value of the mixing ratios. A $0.1 \%$ probability limit separated the acceptable from the unacceptable solutions. In the present case $\chi^{2}$ was equal to $4 \cdot 2$ at the $0 \cdot 1 \%$ limit. Often more than one value of $\chi_{\text {min }}^{2}$ was obtained.

Table 1. Branching ratios for decay of 10 resonances in ${ }^{50} \mathrm{Cr}(\mathrm{p}, \gamma)^{51} \mathrm{Mn}$

| Bound level$E_{\mathrm{x}}(\mathrm{MeV})$ | $E_{\mathrm{p}}=1451$ | 1546 | Branching ratios (\%) from resonances |  |  |  |  |  | 2042 | 2067 keV |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  |  |  | 1580 | 1600 | 1689 | 1798 | 1830 | 2031 |  |  |
| 0 | 8 |  | 5 | 15 | 9 | 34 | 30 | 4 | 18 | 10 |
| 0.238 | 36 |  |  | <1 |  |  | 2 | 11 | 5 |  |
| $1 \cdot 140$ |  |  |  |  |  |  |  |  |  |  |
| 1.817 | 8 |  | 26 |  | 4 | 1 |  | 17 |  |  |
| 1.823 |  | 29 |  | 37 | 7 | 10 | 26 |  | 8 | 5 |
| 1.958 |  | 9 |  |  | 10 | 26 |  | 8 | 7 | 33 |
| 2-139 |  | 13 |  | 23 | 16 | 5 |  | 4 | 9 | 3 |
| $2 \cdot 256$ | 3 |  | 7 |  | 6 |  | 2 |  | 4 | 25 |
| $2 \cdot 275$ |  |  | 15 |  | 3 |  | 3 | 4 |  | 3 |
| $2 \cdot 312$ | 3 |  | 9 |  |  |  | 2 |  | 3 |  |
| $2 \cdot 416$ | 17 |  |  |  |  |  | 1 | 2 |  |  |
| $2 \cdot 702$ | 3 |  | 7 |  |  |  |  | 12 | 3 |  |
| 2.715 |  |  |  |  |  |  |  | 2 |  |  |
| $2 \cdot 846$ |  |  |  |  | 10 | 6 | 4 |  | 8 |  |
| $2 \cdot 892$ | 2 |  |  |  |  |  |  | 2 |  |  |
| 2.914 | 3 | 9 | 2 | $<1$ | 5 | 5 | 1 |  | 1 |  |
| 2.985 |  |  | 3 | 2 | 6 |  | 2 |  | 20 | 11 |
| 3.046 | 6 |  |  |  |  |  |  |  |  |  |
| 3.051 |  |  | 3 |  |  |  | 2 |  |  |  |
| $3 \cdot 131$ | 10 | 3 | 15 |  |  |  | 2 | 9 |  |  |
| 3.294 |  |  | 3 | 3 |  |  | 2 |  |  |  |
| $3 \cdot 423$ |  |  |  |  | 15 |  |  |  |  |  |
| 3.556 |  | 9 |  | 4 | 6 | 5 | 3 |  |  | 2 |
| 3.694 |  |  |  |  |  |  | 3 | 1 |  |  |
| $3 \cdot 896$ | 1 | 26 |  |  |  | 8 | 7 |  | 6 |  |
| 4.052 |  |  |  |  |  |  |  | 1 |  |  |
| 4.203 |  |  |  | 4 |  |  |  | 4 | 5 |  |
| $4 \cdot 355$ |  |  |  | 8 |  |  |  |  |  |  |
| $4 \cdot 453$ |  |  | 5 | 3 |  |  |  |  |  |  |
| 4.729 |  |  |  |  |  |  | 2 |  |  |  |
| $4 \cdot 887$ |  |  |  |  |  |  |  |  | <1 | 2 |
| 4.925 |  |  |  |  |  |  | 3 | 1 |  | <1 |
| 5.073 |  |  |  |  |  |  | 1 | $<1$ | 2 | 1 |
| $5 \cdot 129$ |  |  |  |  |  |  |  | 4 |  |  |
| 5.174 |  |  |  |  |  |  |  | $<1$ |  | <1 |
| 5.188 |  |  |  |  |  |  | <1 | 13 |  |  |
| $5 \cdot 203$ |  |  |  |  | 3 |  |  |  |  |  |
| $5 \cdot 223$ |  | 2 |  |  |  |  |  |  |  | 4 |
| 5.506 |  |  |  |  |  |  |  |  |  | 4 |

## Results

The yield curves found by Wall and Erlandsson (1967) for the reaction ${ }^{50} \mathrm{Cr}(\mathrm{p}, \gamma){ }^{51} \mathrm{Mn}$ were used as a guide to identify the resonances up to $E_{\mathrm{p}}=1830 \mathrm{keV}$. The proton energy of the resonances was obtained using the 992 keV resonance in ${ }^{27} \mathrm{Al}(\mathrm{p}, \gamma)^{28} \mathrm{Si}$ as a calibration point. The yield curve used for resonances above $2 \cdot 0 \mathrm{MeV}$ will be presented in a separate paper.

Table 2. Branching ratios and $\gamma$-ray energies in decay of bound levels in ${ }^{51} \mathbf{M n}$

| Level No. | $\underset{(\mathrm{MeV})}{E_{\mathbf{x}}}$ | Transition | $\begin{gathered} E_{\gamma} \\ (\mathrm{MeV}) \end{gathered}$ | Branching ratio (\%) | Level No. | $\underset{(\mathrm{MeV})}{E_{\mathbf{x}}}$ | Transition | $\begin{gathered} \boldsymbol{E}_{\gamma} \\ (\mathrm{MeV}) \end{gathered}$ | Branching ratio (\%) |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 0 | 0 | g.s. |  |  |  |  |  |  |  |
| 1 | 0.238 | $1 \rightarrow 0$ | $0 \cdot 238$ | 100 | 20 | $3 \cdot 294$ | $20 \rightarrow 0$ | $3 \cdot 294$ | 35 |
| 2 | $1 \cdot 140$ | $2 \rightarrow 0$ | $1 \cdot 140$ | 10 | 21 | $3 \cdot 423$ | $20 \rightarrow 1$ | 3.056 | 65 |
|  |  | $2 \rightarrow 1$ | 0.902 | 90 |  |  | $21 \rightarrow 5$ | $1 \cdot 465$ | 100 |
| 3 | 1.817 | $3 \rightarrow 0$ | $1 \cdot 817$ | 100 | 22 | $3 \cdot 556$ | $22 \rightarrow 0$ | $3 \cdot 556$ | 100 |
| 4 | $1 \cdot 823$ | $4 \rightarrow 0$ | $1 \cdot 823$ | 100 | 23 | $3 \cdot 694$ | $23 \rightarrow 0$$23 \rightarrow 4$$23 \rightarrow 6$ | $\begin{aligned} & 3.694 \\ & 1.871 \\ & 1.555 \end{aligned}$ | $\begin{aligned} & 20 \\ & 55 \\ & 25 \end{aligned}$ |
| 5 | 1.958 | $5 \rightarrow 0$ | 1.958 | 100 |  |  |  |  |  |
|  |  |  |  |  |  |  |  |  |  |
| 6 | 2-139 | $6 \rightarrow 0$ | 2-139 | 100 | 24 | $3 \cdot 896$ | $24 \rightarrow 0$ | $3 \cdot 896$ | 100 |
| 7 | $2 \cdot 256$ | $7 \rightarrow 0$ | $2 \cdot 256$ | 20 | 25 | $4 \cdot 052$ | $25 \rightarrow 16$ | $1 \cdot 067$ | 100 |
|  |  | $7 \rightarrow 1$ | $2.018$ | $80$ |  |  |  |  |  |
| 8 | $2 \cdot 275$ |  |  |  | 26 | 4-203 | $26 \rightarrow 0$ | $4 \cdot 203$ | 20 |
|  |  | $8 \rightarrow 3$ | 0.458 | 70 |  |  | $26 \rightarrow 1$ | $3.965$ | 20 |
|  |  | $8 \rightarrow 5$ | $0 \cdot 317$ | $30$ |  |  | $26 \rightarrow 4$ | $2 \cdot 380$ | 60 |
| 9 | $2 \cdot 312$ | $9 \rightarrow 0$ | $2 \cdot 312$ | 15 | 27 | $4 \cdot 355$ | $27 \rightarrow 0$ | $4 \cdot 355$ | 452515 |
|  |  | $9 \rightarrow 1$ | $2 \cdot 074$ | 85 |  |  | $27 \rightarrow 5$ | $2 \cdot 397$ |  |
| 10 | $2 \cdot 416$ | $10 \rightarrow 0$ | $2 \cdot 416$ | 6 |  | $4 \cdot 453$ | $27 \rightarrow 22$ | 0.799 |  |
|  |  | $10 \rightarrow 1$ | $2 \cdot 178$ | 34 | 28 |  | $28 \rightarrow 4$$28 \rightarrow 5$$28 \rightarrow 9$ | $2 \cdot 630$$2 \cdot 495$$2 \cdot 141$ | 452035 |
|  |  | $10 \rightarrow 2$ | 1.276 | 37 |  |  |  |  |  |
|  |  | $10 \rightarrow 3$ | $0 \cdot 599$ | 15 |  |  |  |  |  |
|  |  | $10 \rightarrow 6$ | $0 \cdot 277$ | 8 |  |  |  |  |  |
| 11 | $2 \cdot 702$ | $11 \rightarrow 2$ | 1.562 | 2 | 29 | $4 \cdot 729$ | $29 \rightarrow 0$ | 4•729 | 85 |
|  |  | $11 \rightarrow 3$ | 0.885 | 21 | 30 | 4-887 | $30 \rightarrow 5$ |  |  |
|  |  | $11 \rightarrow 5$ | 0.744 | 26 |  |  | $30 \rightarrow 13$ | $2 \cdot 041$ | 60 |
|  |  | $11 \rightarrow 9$ | $0 \cdot 390$ | 37 |  |  |  |  |  |
| 12 | $2 \cdot 715$ | $12 \rightarrow 0$ | $2 \cdot 715$ | 100 | 31 | 4.925 | $31 \rightarrow 15$ | $2 \cdot 011$ | 75 |
|  |  |  |  |  |  |  | $31 \rightarrow 16$ | $1.940$ | 25 |
| 13 | $2 \cdot 846$ | $13 \rightarrow 4$ | 1.023 | 40 | 32 | $5 \cdot 073$ | $32 \rightarrow 3$ | 3-256 |  |
|  |  | $13 \rightarrow 5$ | 0.888 | 60 |  |  | $32 \rightarrow 5$ | 3.115 | 35 |
| 14 | $2 \cdot 892$ | $14 \rightarrow 0$ | 2.892 | 35 | 33 | 5•129 | $32 \rightarrow 8$ | $2 \cdot 798$ | 50 |
|  |  | $14 \rightarrow 2$ | 1.752 | 45 |  |  | $33 \rightarrow 9$ | $2 \cdot 817$ | 75 |
|  |  | $14 \rightarrow 3$ | 1.075 | 20 |  |  | $33 \rightarrow 16$ | 2.144 | 75 25 |
| 15 | 2.914 | $15 \rightarrow 0$ | $2 \cdot 914$ | 100 | 34 | 5•174 | 34 $\rightarrow$ 4 | $3 \cdot 351$ | 100 |
| 16 | $2 \cdot 985$ | $\begin{aligned} & 16 \rightarrow 0 \\ & 16 \rightarrow 1 \end{aligned}$ | $\begin{aligned} & 2.985 \\ & 2.747 \end{aligned}$ | 20 | 35 | $5 \cdot 188$ | $\begin{aligned} & 35 \rightarrow 1 \\ & 35 \rightarrow 14 \\ & 35 \rightarrow 15 \\ & 35 \rightarrow 18 \end{aligned}$ | $\begin{aligned} & 4 \cdot 950 \\ & 2 \cdot 296 \end{aligned}$ | $\begin{aligned} & 34 \\ & 18 \end{aligned}$ |
|  |  |  |  | 80 |  |  |  |  |  |
| 17 | $3 \cdot 046$ | $\begin{aligned} & 17 \rightarrow 3 \\ & 17 \rightarrow 7 \\ & 17 \rightarrow 8 \end{aligned}$ | $\begin{aligned} & 1.229 \\ & 0.790 \\ & 0.771 \end{aligned}$ | 40 <br> 45 <br> 15 |  |  |  | 2.274 | 20 |
|  |  |  |  |  |  |  |  | 2-137 | 28 |
|  |  |  |  |  | 36 | $5 \cdot 203$ | $36 \rightarrow 0$ | $5 \cdot 203$ | 70 |
| 18 | $3 \cdot 051$ | $18 \rightarrow 0$ | $3 \cdot 051$ | 100 | 37 | $5 \cdot 223$ | $37 \rightarrow 1$ | 4.985 | 50 |
| 19 | $3 \cdot 131$ | $19 \rightarrow 3$ | 1.314 | 45 | 38 |  | $37 \rightarrow 3$ | $3 \cdot 406$ | 30 |
|  |  | $19 \rightarrow 7$ | 0.875 | 12 |  | $5 \cdot 506$ | $\begin{aligned} & 38 \rightarrow 1 \\ & 38 \rightarrow 25 \end{aligned}$ | $\begin{aligned} & 5 \cdot 268 \\ & 1 \cdot 454 \end{aligned}$ |  |
|  |  | $19 \rightarrow 9$ | 0.819 | 10 |  |  |  |  | $\begin{aligned} & 35 \\ & 65 \end{aligned}$ |

Figs 1a-1f (pp. 267-72). Gamma-ray spectra measured at $55^{\circ}$ with a $35 \mathrm{~cm}^{3} \mathrm{Ge}(\mathrm{Li})$ detector for resonances at $E_{\mathrm{p}}=1689,1798,1830,2031,2042$ and 2067 keV in the reaction ${ }^{50} \mathrm{Cr}(\mathrm{p}, \gamma)^{51} \mathrm{Mn}$.








Fig. 2. Decay schemes of 10 resonance states ( $E_{\mathrm{p}}=1451-2067 \mathrm{keV}$ ) to bound levels in ${ }^{51} \mathrm{Mn}$.

## Resonance decay schemes

The resonances at $E_{\mathrm{p}}=1451,1546,1580,1600,1798$ and 1830 keV have already been investigated by previous workers, but several revisions in the decay schemes are found to be necessary. The resonances at $E_{\mathrm{p}}=1689,2031,2042$ and 2067 keV have not been studied earlier. The decay of the resonance states is summarized in Table 1, while the branching ratios for the transitions in the decay of the bound states and the $\gamma$-ray energies are given in Table 2.
1451 keV Resonance ( $E_{\mathrm{x}}=6.689 \mathrm{MeV}$ )
The decay scheme found for this resonance is in agreement with the results of Forsblom et al. (1972), though the branching ratios differ for certain transitions. The weak transition to the level at 3.896 MeV was not reported by them.
1546 keV Resonance ( $E_{\mathrm{x}}=6.785 \mathrm{MeV}$ )
Again the decay scheme of this resonance agrees in general with the results reported by Forsblom et al. (1972). However, their reported $2 \%$ transition to the ground state and $1 \%$ transition to the bound level at $5 \cdot 585 \mathrm{MeV}$ were not detected in our study. The $3 \%$ transition to the level at $3 \cdot 131 \mathrm{MeV}$ observed here was not reported by them. 1580 keV Resonance ( $E_{\mathrm{x}}=6.818 \mathrm{MeV}$ )

The transitions observed from this resonance level to the bound levels at $2 \cdot 985$, 3.051 and 3.294 MeV were not reported by Forsblom et al. (1972), while transitions noted by them to the levels at 2.841 and 3.422 MeV were not seen in our $\gamma$-ray spectra. The rest of the decay scheme is in agreement.
1600 keV Resonance ( $E_{\mathrm{x}}=6.838 \mathrm{MeV}$ )
The decay scheme of this resonance is in general agreement with the results of Wall and Erlandsson (1967) and Forsblom et al. (1972). Those authors did not record, however, the transitions from the resonance to the first excited state and to the level at 4.453 MeV . Their reported transitions to $2 \cdot 276,2 \cdot 488$ and 3.058 MeV levels were not observed in the present investigation. The population of the level at 4.453 MeV from the two resonances at 1580 and 1600 keV confirms its existence. 1689 keV Resonance ( $E_{\mathrm{x}}=6.926 \mathrm{MeV}$ )

The resonance at 1689 keV was reported by Wall and Erlandsson (1967) and Forsblom et al. (1972), but decay schemes were not determined by them. The $\gamma$-ray spectrum from the present work is shown in Fig. 1a. The resonance decays to a number of bound levels as shown in the decay scheme (Fig. 2). The population of the bound level at 5.203 MeV was observed for the first time in the proton capture reaction.

## 1798 keV Resonance ( $E_{\mathrm{x}}=7.032 \mathrm{MeV}$ )

The results for the decay of the 1798 keV resonance (Fig. 1b) agree in general with those of Wall and Erlandsson (1967). However, it is found that the resonance has a $6 \%$ decay to the 2.846 MeV level which was not reported in the previous work.
1830 keV Resonance ( $E_{\mathrm{x}}=7.062 \mathrm{MeV}$ )
The $\gamma$-ray spectrum measured at this resonance is shown in Fig. $1 c$; the decay scheme resulting from the analysis of the spectrum is presented in Fig. 2. The results obtained by Wall and Erlandsson (1967) are considerably different from the present findings. Here transitions were detected from the resonance to the first excited state
and to the levels at $2 \cdot 275,2 \cdot 416,2 \cdot 846,2 \cdot 985,3 \cdot 051,3 \cdot 556,4 \cdot 729,4 \cdot 925,5 \cdot 073$ and $5 \cdot 188 \mathrm{MeV}$, all of which were not found in the previous work. A transition from the resonance to the level at 4.017 MeV as reported by Wall and Erlandsson was not observed in the present study. The resonance branching ratios given here (Fig. 2) are also different from their values. The levels at $4 \cdot 729,4 \cdot 925,5 \cdot 073$ and $5 \cdot 188 \mathrm{MeV}$ have been observed to be populated in the proton capture reaction for the first time.

## 2031 keV Resonance $\left(E_{\mathrm{x}}=7 \cdot 259 \mathrm{MeV}\right.$ )

The $\gamma$-ray spectrum for this resonance, which has not been previously investigated, is shown in Fig. 1d, the decay scheme resulting from the analysis of the spectrum being incorporated in Fig. 2. The resonance level decays to a number of bound levels, with relatively stronger ( $10 \%$ ) transitions to the levels at $0 \cdot 238,1 \cdot 817,2 \cdot 702$ and $5 \cdot 188 \mathrm{MeV}$. The population of the levels at $4.925,5 \cdot 073$ and $5 \cdot 188 \mathrm{MeV}$ confirms their existence, as these levels were also observed at the $E_{\mathrm{p}}=1830 \mathrm{keV}$ resonance. The bound levels at $4 \cdot 052,5 \cdot 129$ and $5 \cdot 174 \mathrm{MeV}$ have been observed to be populated for the first time in the proton capture reaction.
2042 keV Resonance ( $E_{\mathrm{x}}=7.270 \mathrm{MeV}$ )
The $\gamma$-ray spectrum obtained at $E_{\mathrm{p}}=2042 \mathrm{keV}$ is shown in Fig. $1 e$ and the resulting decay scheme is included in Fig. 2. This resonance has not been investigated previously. The resonance state decays to 15 bound levels including the ground state. The strongest transitions with $20 \%$ and $18 \%$ branches are to the bound level at 2.985 MeV and the ground state respectively.

2067 keV Resonance ( $E_{\mathrm{x}}=7.295 \mathrm{MeV}$ )
The $\gamma$-ray spectrum obtained at the 2067 keV resonance is shown in Fig. $1 f$ and the resulting decay scheme is included in Fig. 2. This resonance has not been investigated previously. The strong transitions exhibited by the resonance are $11 \%$, $25 \%, 33 \%$ and $10 \%$ to the bound levels at $2 \cdot 985,2 \cdot 256,1 \cdot 958 \mathrm{MeV}$ and the ground state respectively. There are, however, weaker transitions to a number of bound levels at $4 \cdot 887,5 \cdot 073$ and $5 \cdot 506 \mathrm{MeV}$ among others.

## Properties of bound levels

The decay schemes of the bound levels as obtained from the study of the resonances are given in Fig. 3. In the following, some comments are made for either those levels where the decay schemes differ from earlier reports or the levels which have been observed for the first time in the proton capture reaction.

## 1. 140 MeV Level

The level at $1 \cdot 140 \mathrm{MeV}$ was not populated directly from any of the resonances. Its existence is based only on the decays of the bound levels at $2 \cdot 416,2 \cdot 702$ and $2 \cdot 892 \mathrm{MeV}$. The level decays $10 \%$ to the ground state and $90 \%$ to the first excited state, which is in agreement with the results of Forsblom et al. (1972).
Doublet Levels at 1.817 and 1.823 MeV
Both of these levels decay completely to the ground state. Their existence is confirmed by the presence of the primary $\gamma$-rays from different resonances. The strongest transitions $(26 \%$ and $17 \%)$ to the level at 1.817 MeV are from the
resonances at 1580 and 2031 keV , whereas $29 \%$ and $26 \%$ decays to the level at 1.823 MeV come from two resonances at 1546 and 1830 keV .

### 2.702 and 2.715 MeV Levels

The decay schemes of both these levels were derived from the resonance at 2031 keV , from which the 2.702 MeV level was populated relatively strongly compared with other resonances. This level has $2 \%, 21 \%, 26 \%$ and $37 \%$ decays to the levels at $1 \cdot 140,1 \cdot 817,1 \cdot 958$ and $2 \cdot 312 \mathrm{MeV}$ respectively. The results disagree with those of Forsblom et al. (1972), who reported the decays only to the levels at 1.958 and 2.312 MeV . The bound level at 2.715 MeV decays completely to the ground state, rather than to the level at 1.817 MeV as reported by Wall and Erlandsson (1967).


Fig. 3. Decay schemes of low lying bound levels of ${ }^{51} \mathrm{Mn}$ up to an excitation energy of 5.51 MeV , as determined in the present investigation.

### 2.846 MeV Level

The transition ratios of $40 \%$ and $60 \%$ from the bound level at $2 \cdot 846 \mathrm{MeV}$ to the levels at 1.823 and 1.958 MeV are in disagreement with the findings of Forsblom et al. (1972), who reported $25 \%$ and $75 \%$ transitions respectively.

## $3 \cdot 046 \mathrm{MeV}$ Level

This bound level showed three transitions, $40 \%, 45 \%$ and $15 \%$ leading to the levels at $1 \cdot 817,2 \cdot 256$ and 2.275 MeV respectively. The transition to the first excited state reported by Forsblom et al. (1972) was not observed.

## 3•294 MeV Level

The bound level at $3 \cdot 294 \mathrm{MeV}$ has only two transitions: $35 \%$ to the ground state and $65 \%$ to the first excited state. These values are to be compared with $25 \%$ and $75 \%$ as reported by Forsblom et al. (1972).

### 3.694 MeV Level

The bound level at 3.694 MeV was populated in the 1830 keV resonance. It decays $20 \%, 55 \%$ and $25 \%$ to the ground state and the levels at $1 \cdot 823$ and $2 \cdot 139$ MeV respectively. The $10 \%$ decay to the level at 2.914 MeV reported by Forsblom et al. (1972) was not observed.

### 3.896 MeV Level

The strong transition ( $26 \%$ ) to the bound level at 3.896 MeV was observed in the 1546 keV resonance. The level decays $100 \%$ to the ground state. The transition to the level at $2 \cdot 846 \mathrm{MeV}$ as reported by Forsblom et al. (1972) was not observed.

## $4 \cdot 052 \mathrm{MeV}$ Level

The bound level at 4.052 MeV was populated (weakly) at one resonance, namely $E_{\mathrm{p}}=2031 \mathrm{keV}$. The level has not been reported earlier and it decays $100 \%$ to the level at 2.985 MeV .

### 4.453 MeV Level

This level was excited at two resonances, 1580 and 1600 keV . It decays $45 \%$, $20 \%$ and $35 \%$ to the levels at $1 \cdot 823,1.958$ and $2 \cdot 312 \mathrm{MeV}$; Forsblom et al. (1972) only reported a $46 \%$ decay to the level at 1.823 MeV . The bound level at 4.453 MeV has been identified as an isobaric analogue associated with the ground state of ${ }^{51} \mathrm{Cr}$ by Rapaport et al. (1967).

### 4.729 and 4.925 MeV Levels

These levels are excited at the 1830 keV resonance. They were observed to be populated for the first time in the proton capture reaction, though they were reported in ( ${ }^{3} \mathrm{He}, \mathrm{d}$ ) studies by Rapaport et al. (1967). Only $85 \%$ of the decay from the level at 4.729 MeV to the ground state is accounted for by the present results. The level at 4.925 MeV decays $75 \%$ and $25 \%$ to the levels at 2.914 and 2.985 MeV respectively.

### 4.887 MeV Level

This level was excited at two resonances, 2042 and 2067 keV . It decays $60 \%$ to the level at 2.846 MeV and $13 \%$ to the level at 1.958 MeV . The remaining $27 \%$ of its decay could not be determined from the $\gamma$-ray spectra.

## $5 \cdot 073 \mathrm{MeV}$ Level

This level was reported by Rapaport et al. (1967) from the ( ${ }^{3} \mathrm{He}, \mathrm{d}$ ) reaction, but it has been observed to be populated in the proton capure reaction here for the first time. The level decays to the $2.275,1.958$ and 1.817 MeV levels with branching ratios of $50 \%, 35 \%$ and $15 \%$ respectively.

## $5 \cdot 129 \mathrm{MeV}$ Level

This level was also seen to be excited for the first time in the proton capture reaction, though it was observed in the ( ${ }^{3} \mathrm{He}, \mathrm{d}$ ) reaction by Rapaport et al. (1967). It decays $75 \%$ to the level at $2 \cdot 312 \mathrm{MeV}$ and $25 \%$ to the level at 2.985 MeV .

## $5 \cdot 174 \mathrm{MeV}$ Level

Again this level was seen to be excited for the first time in the proton capture reaction, though it was reported in the ( $\left.{ }^{3} \mathrm{He}, \mathrm{d}\right)$ reaction by Rapaport et al. (1967). It decays $100 \%$ to the level at 1.823 MeV with a $\gamma$-ray of energy 3.35 MeV .

### 5.188 MeV Level

This level decays to the first excited state with the emission of a $4.95 \mathrm{MeV} \gamma$-ray $(34 \%)$, to the 2.892 MeV level with a $2.297 \mathrm{MeV} \gamma$-ray $(18 \%)$, to the 2.914 MeV level with a $2.274 \mathrm{MeV} \gamma$-ray $(20 \%)$ and to the 3.051 MeV level with a 2.137 MeV $\gamma$-ray $(28 \%)$. The level has not been observed in charged particle reactions.

### 5.203 MeV Level

This level was only excited at $E_{\mathrm{p}}=1689 \mathrm{keV}$. A $70 \%$ decay goes to the ground state.

### 5.506 MeV Level

This level was reported by Rapaport et al. (1967), but again it is observed in the proton capture reaction for the first time here. Its decay is to two levels: $35 \%$ to the first excited state and $65 \%$ to the level at 4.052 MeV .

Table 3. $\boldsymbol{\gamma}$-ray angular distribution measurements at two resonances in ${ }^{51} \mathbf{M n}$ The results summarized are the Legendre coefficients $A_{2}=a_{2} / a_{0}$ and $A_{4}=a_{4} / a_{0}$ from

$$
W(\theta)=a_{0}+a_{2} P_{2}(\theta)+a_{4} P_{4}(\theta)
$$

the values of $\chi_{\text {min }}^{2}$ and the mixing ratios $\delta$ for the primary $\gamma$-rays emitted in the transitions from the resonances to the final states $E_{f}$

| Final state |  | Experimental |  | Theoretical |  | $\chi_{\text {min }}^{2}$ | $\delta$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| $E_{\mathrm{f}}(\mathrm{MeV})$ | $J_{\text {f }}$ | $A_{2}$, | $A_{4}$ | $A_{2}$ | $A_{4}$ |  |  |
| $E_{\mathrm{p}}=2042 \mathrm{keV}\left(E_{\mathrm{x}}=7 \cdot 270 \mathrm{MeV}, \mathrm{J}=5 / 2\right)$ |  |  |  |  |  |  |  |
| 0 | 5/2 | $-0.20 \pm 0.04$ | $-0.11 \pm 0.04$ | $-0 \cdot 19$ | $-0 \cdot 12$ | 0.43 | $0 \cdot 60 \pm 0 \cdot 02$ |
| 0.238 | 7/2 | $-0.02 \pm 0.05$ | $-0.16 \pm 0.05$ | $-0.02$ | $-0 \cdot 16$ | $\begin{aligned} & 11 \cdot 8 \\ & 8 \cdot 90 \end{aligned}$ | $\begin{gathered} -0 \cdot 18 \pm 0 \cdot 08 \\ 5 \cdot 67_{-2.2}^{+8.6} \end{gathered}$ |
| $1 \cdot 823$ | 3/2 | $0 \cdot 15 \pm 0.05$ | $0 \cdot 13 \pm 0 \cdot 06$ | $0 \cdot 16$ | $0 \cdot 13$ | $3 \cdot 30$ | $-0.29 \pm 0.05$ |
| $2 \cdot 139$ | 3/2 | $0.03 \pm 0.06$ | $-0.11 \pm 0.06$ | 0.04 | $-0.12$ | 2.86 | $-0.30 \pm 0.07$ |
| 2.985 | 5/2 | $-0.02 \pm 0.05$ | $-0.21 \pm 0.05$ | $-0.02$ | $-0.21$ | $2 \cdot 80$ | $0 \cdot 49 \pm 0 \cdot 11$ |
|  |  |  |  |  |  | $0 \cdot 28$ | $-6 \cdot 30 \pm 2 \cdot 3$ |
| $E_{\mathrm{p}}=2067 \mathrm{keV}\left(E_{\mathrm{x}}=7 \cdot 295 \mathrm{MeV}, J=5 / 2\right)$ |  |  |  |  |  |  |  |
| 0 | 5/2 | $0.44 \pm 0.06$ | $-0.13 \pm 0.06$ | $0 \cdot 45$ | $-0 \cdot 13$ | $1 \cdot 22$ | $-1.43 \pm 0.36$ |
| 1.958 | 3/2 | $-0.18 \pm 0.03$ | $-0 \cdot 10 \pm 0.03$ | $-0 \cdot 18$ | $-0 \cdot 10$ | $3 \cdot 31$ | $0.27 \pm 0.05$ |
|  |  |  |  |  |  | $2 \cdot 44$ | $4 \cdot 33 \pm 0.03$ |
| $2 \cdot 256$ | 7/2 | $-0 \cdot 02 \pm 0 \cdot 04$ | $-0 \cdot 10 \pm 0 \cdot 04$ | -0.03 | $-0 \cdot 10$ | 3.90 | $-0 \cdot 10 \pm 0 \cdot 19$ |
|  |  |  |  |  |  | $1 \cdot 86$ | $5 \cdot 14_{-2.54}^{+4 \cdot 36}$ |
| $2 \cdot 985$ | 5/2 | $-0.06 \pm 0.03$ | $-0 \cdot 20 \pm 0.03$ | -0.06 | $-0 \cdot 20$ | $2 \cdot 20$ | $-8 \cdot 14_{-6.16}^{+2.54}$ |
|  |  |  |  |  |  | $9 \cdot 00$ | $0 \cdot 53 \pm 0.03$ |

## Spins of resonance levels

The results of the angular distribution measurements for the resonances at $E_{\mathrm{p}}=2042$ and 2067 keV are summarized in Table 3, which gives the Legendre
coefficients, the transition mixing ratios for the primary $\gamma$-rays from ${ }^{51} \mathrm{Mn}$ resonances and the values of $\chi_{\text {min }}^{2}$.
2042 keV Resonance ( $E_{\mathrm{x}}=7 \cdot 270 \mathrm{MeV}$ )
Since the measured distributions of the five transitions from the resonance level at $E_{\mathrm{x}}=7 \cdot 270 \mathrm{MeV}$ all show clear anisotropy, a spin of $1 / 2$ can be excluded. The data and the chi-squared projection curves are shown in Fig. 4. It is clear from the figure that a spin of $5 / 2$ for the resonance state has the minimum $\chi^{2}$ value. The choice of $5 / 2$ is also consistent with the transitions to the bound levels at $0.238 \mathrm{MeV}(7 / 2)$, $1 \cdot 823 \mathrm{MeV}(3 / 2), 2 \cdot 139 \mathrm{MeV}(3 / 2)$ and $2 \cdot 985 \mathrm{MeV}(5 / 2)$. The $\chi_{\text {min }}^{2}$ values are all below the $0 \cdot 1 \%$ limit except for the transition to the first excited state.


Fig. 4. Results of the analysis of the angular distributions measured for the $\gamma$-ray transition from the $E_{\mathrm{p}}=2042 \mathrm{keV}$ resonance to the ground state.

2067 keV Resonance $\left(E_{\mathrm{x}}=7 \cdot 295 \mathrm{MeV}\right.$ )
As before, the measured distributions of the four transitions from the resonance level at $E_{\mathrm{x}}=7 \cdot 295 \mathrm{MeV}$ show clear anisotropy, and a spin of $1 / 2$ is thus excluded. The $\chi_{\min }^{2}$ value for the transition to the ground state is within the $0 \cdot 1 \%$ acceptance limit if the spin of the resonance state is $5 / 2$. For the other three transitions to the levels at $1.958 \mathrm{MeV}(3 / 2), 2 \cdot 256 \mathrm{MeV}(7 / 2)$ and $2 \cdot 985 \mathrm{MeV}(5 / 2)$, the $\chi_{\text {min }}^{2}$ is again within the limit for the same spin assignment. Therefore, a spin of $5 / 2$ for the resonance at $E_{\mathrm{p}}=2067 \mathrm{keV}$ is favoured (Fig. 5).


Fig. 5. Results of the analyses of the angular distributions measured for the $\gamma$-ray transitions from the $E_{\mathrm{p}}=2067 \mathrm{keV}$ resonance to the two bound levels at 1.958 and 2.985 MeV .

## Discussion

The study of 10 resonances for the reaction ${ }^{50} \mathrm{Cr}(\mathrm{p}, \gamma){ }^{51} \mathrm{Mn}$ has provided extensive information on the energy levels, the spins of two of the resonance states and the decay properties of the bound levels in ${ }^{51} \mathrm{Mn}$. Table 4 compares the excitation energies of ${ }^{51} \mathrm{Mn}$ as found from investigations of $(\mathrm{p}, \gamma)$ and $\left({ }^{3} \mathrm{He}, \mathrm{d}\right)$ reactions, including the present study. There are 38 bound levels up to an excitation of 5.506 MeV which have been populated from the 10 resonances. The two resonances at 1830

Table 4. Excitation energies of ${ }^{51} \mathrm{Mn}$ bound levels up to $5 \cdot 506 \mathrm{MeV}$
The present results from the ${ }^{50} \mathrm{Cr}(\mathrm{p}, \gamma)^{51} \mathrm{Mn}$ reaction are compared with those of (F) Forsblom et al. (1972) and (W) Wall and Erlandsson (1967) from the same reaction and with those of (R) Rapaport et al. (1967) from the ${ }^{50} \mathrm{Cr}\left({ }^{3} \mathrm{He}, \mathrm{d}\right){ }^{51} \mathrm{Mn}$ reaction

|  | $\begin{gathered} \text { F } \\ (\mathrm{p}, \gamma) \\ (\mathrm{MeV}) \end{gathered}$ |  |  | $\begin{aligned} & \text { Present } \\ & (\mathrm{p}, \gamma) \\ & (\mathrm{MeV}) \end{aligned}$ | $\begin{gathered} \mathrm{F} \\ (\mathrm{p}, \gamma) \\ (\mathrm{MeV}) \end{gathered}$ | $\begin{gathered} \text { W } \\ (\mathrm{p}, \gamma) \\ (\mathrm{MeV}) \end{gathered}$ | $R$ $\left({ }^{3} \mathrm{He}, \mathrm{d}\right)$ $(\mathrm{MeV})$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| $0 \cdot 238$ | 0.238 | $0 \cdot 240$ | 0.240 | $3 \cdot 423$ | 3.423 | $3 \cdot 422$ | $3 \cdot 427$ |
| $1 \cdot 140$ | $1 \cdot 140$ |  | $1 \cdot 160$ | $3 \cdot 556$ | $3 \cdot 553$ | $3 \cdot 549$ | $3 \cdot 555$ |
|  |  |  | $1 \cdot 500$ | $3 \cdot 694$ | 3.694 | $3 \cdot 690$ | $3 \cdot 698$ |
| 1.817 | 1.817 | 1.817 |  | $3 \cdot 896$ | 3•894 | $3 \cdot 897$ | 3.900 |
| 1.823 | 1.825 | $1 \cdot 826$ | 1.830 |  |  | $4 \cdot 017$ | $4 \cdot 017$ |
| 1.958 | 1.959 | 1.958 | 1.962 | $4 \cdot 052$ |  |  |  |
| 2.139 | 2-140 | $2 \cdot 142$ | $2 \cdot 147$ |  | 4.091 |  | $4 \cdot 099$ |
| $2 \cdot 256$ | $2 \cdot 256$ | $2 \cdot 257$ |  | $4 \cdot 203$ | $4 \cdot 206$ |  |  |
| 2. 275 | $2 \cdot 276$ | $2 \cdot 276$ | $2 \cdot 284$ | $4 \cdot 355$ | $4 \cdot 352$ |  | $4 \cdot 367$ |
| $2 \cdot 312$ | $2 \cdot 310$ | $2 \cdot 311$ |  | $4 \cdot 453$ | $4 \cdot 451$ |  | $4 \cdot 446$ |
| $2 \cdot 416$ | $2 \cdot 416$ |  | $2 \cdot 426$ |  | $4 \cdot 488$ |  | $4 \cdot 497$ |
|  |  | $2 \cdot 488$ |  |  | $4 \cdot 540$ |  | $4 \cdot 533$ |
|  |  | $2 \cdot 660$ |  | $4 \cdot 729$ |  |  | 4•730 |
| $2 \cdot 702$ | $2 \cdot 702$ | $2 \cdot 702$ |  | $4 \cdot 887$ | $4 \cdot 883$ |  |  |
| $2 \cdot 715$ |  | $2 \cdot 717$ |  | $4 \cdot 925$ |  |  | $4 \cdot 925$ |
|  |  | 2.808 |  | $5 \cdot 073$ |  |  | $5 \cdot 074$ |
| $2 \cdot 846$ | $2 \cdot 841$ |  | $2 \cdot 844$ | 5.129 |  |  | $5 \cdot 128$ |
| $2 \cdot 892$ | $2 \cdot 893$ |  |  | $5 \cdot 174$ |  |  | 5•176 |
| $2 \cdot 914$ | 2.914 | $2 \cdot 914$ | $2 \cdot 919$ | $5 \cdot 188$ |  |  |  |
| $2 \cdot 985$ | $2 \cdot 985$ | $2 \cdot 981$ | 2-985 | $5 \cdot 203$ |  |  |  |
| $3 \cdot 046$ | $3 \cdot 049$ |  |  |  | $5 \cdot 212$ |  |  |
| $3 \cdot 051$ | 3.058 |  | $3 \cdot 058$ | $5 \cdot 223$ | $5 \cdot 223$ |  | $5 \cdot 225$ 5.444 |
|  | $3 \cdot 091$ | 3.094 |  |  |  |  | $5 \cdot 444$ 5.508 |
| 3•131 | $3 \cdot 130$ | 3.128 | 3-135 | 5•506 |  |  | 5.508 |
| 3-294 | 3-292 | 3.285 | $3 \cdot 300$ |  |  |  |  |

and 2031 keV alone populated 20 and 19 bound levels respectively. Nine bound levels at $4 \cdot 052,4 \cdot 729,4 \cdot 925,5 \cdot 073,5 \cdot 129,5 \cdot 174,5 \cdot 188,5 \cdot 203$ and $5 \cdot 506 \mathrm{MeV}$ have been observed for the first time in the proton capture reaction, including the three levels at $4 \cdot 052,5 \cdot 188$ and $5 \cdot 203 \mathrm{MeV}$ which have not been detected in charged particle reactions either. The previously unknown decay properties of the bound levels at $4 \cdot 729,4 \cdot 925,5 \cdot 073,5 \cdot 129,5 \cdot 174,5 \cdot 188$ and $5 \cdot 506 \mathrm{MeV}$ have been determined, while revisions in the decay schemes have been found to be necessary for the levels at $2 \cdot 416,2 \cdot 702,2 \cdot 715,3 \cdot 046,3 \cdot 294,3 \cdot 423,3 \cdot 694,3 \cdot 896,4 \cdot 355$, 4.453 and 5.223 MeV . The derived values of level energies generally differ from
one reaction to another, possibly due to the differences in energy resolution of the particular detecting systems. A $Q$ value to the ground state of $5 \cdot 269 \pm 0.002 \mathrm{MeV}$ for the reaction ${ }^{50} \mathrm{Cr}(\mathrm{p}, \gamma){ }^{51} \mathrm{Mn}$ has been derived from the known proton bombarding energies and the energies of the strong direct ground-state transitions from the resonance levels, as determined from the $\gamma$-ray spectra observed with the $\mathrm{Ge}(\mathrm{Li})$ detector.

The interesting point in the present data is that the second excited state at $1 \cdot 140$ MeV was not populated directly in any of the resonances studied. The same is true in the results of Wall and Erlandsson (1967) and Forsblom et al. (1972). The level was only populated through the cascades from the bound levels at $2 \cdot 416,2 \cdot 702$ and $2 \cdot 892 \mathrm{MeV}$. One of the reasons for this could be that the level has a spin of $\geqslant 9 / 2$, and resonances studied below 2 MeV proton bombarding energies are less likely to have higher spin values because of penetrability considerations. At higher proton energies it might be possible to form resonances of higher spins which would populate the level at $1 \cdot 140 \mathrm{MeV}$ and other high spin bound levels as well. There have been theoretical predictions of states with spin $9 / 2$ and $11 / 2$ below 3 MeV excitation energies in ${ }^{51} \mathrm{Mn}$ by Malik and Scholz (1966) and McCullen et al. (1964). The probable values of the bound levels (not shown in Figs 2 and 3) as reported by Forsblom et al. (1972) indicate only one $9 / 2$ spin, for the level at $1 \cdot 140 \mathrm{MeV}$. This is an added indication that higher spin resonances are needed to populate bound levels with spin $\geqslant 9 / 2$. Further work at higher proton bombarding energies might help to check the theoretical predictions.

## Acknowledgment

One of the authors (G.U.D.) would like to thank Professor E. W. Titterton for the encouragement to carry out the research work at the Department of Nuclear Physics, Australian National University.

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